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Refinement of plume models in
the infrared spectral region

By

Kenneth E. Harwell

University of Tennessee Space Institute
30 June 1978

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The MIRADCOM Radiation Model was improved and used successfully to predict CO ₂ and H ₂ O radiation emitted from a turbojet exhaust plume and a rocket exhaust plume. The model was validated using experimental data. Two existing gas dynamics models (LAPP and REP3) were employed to calculate the flow field properties required as input to the radiation computer codes. A new finite difference code was developed, but was only partially successfully in predicting imbedded shock waves.		

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Final Report for
Contract DAAK40-77-C-0032

REFINEMENT OF PLUME MODELING IN
THE INFRARED SPECTRAL REGION

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The University of Tennessee Space Institute
Tullahoma, Tennessee 37388

30 June 1978

Approved for public release; distribution unlimited

Prepared for
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ABSTRACT

A research program was conducted to develop a validated numerical model for the prediction of radiation emitted in the infrared wavelength spectral band from turbojet and rocket exhaust plumes. Two existing gas dynamics models (LAPP and REP3) were employed to calculate the flow field properties (temperature and species concentrations) required as input for the radiation computer codes. A new finite difference computer code, which was developed as part of this project, and includes the calculation of imbedded shock waves in the exhaust plume along with finite-rate chemistry and turbulent mixing, was only partially successful. A refined version of the MIRADCOM Radiation Model was developed and was used successfully to predict CO₂ and H₂O radiation emitted from several cases for which experimental data were available.

1. INTRODUCTION

During the past few years, considerable effort has been expended to design infrared seeker systems for ground-to-air and air-to-air missile guidance systems. Since the next generation of seeker systems must possess head-on capability, infrared radiation from the missile exhaust plume is an important component of the overall infrared signature of a missile system. In order to characterize completely the missile seeker-target interaction, it is necessary to develop infrared radiation predictive models which have the capability of predicting all the pertinent physical phenomena associated with the plume infrared signature.

Infrared radiation emitted from the hot gases in an exhaust plume is characterized by the types of radiating species and their number densities and the environmental conditions (temperature and pressure). All radiative predictive models require a knowledge of these parameters, so that the first step in developing a model is to develop or to obtain a gas dynamics model which yields the spatial distribution of radiating specie concentrations, temperature, and pressure.

Unfortunately, the exhaust plumes of turbojet aircraft and rockets are characterized by the existence of a complex shock wave structure and the presence of a turbulent viscous mixing layer between the inner hot jet flow and the external freestream flow. The inner flow usually consists of a high-temperature, high-velocity, chemically-reacting gas produced by the combustion of a hydrocarbon fuel while the external flow usually is low-temperature, low-velocity air that flows around the exterior of the aircraft or missile body.

Even though considerable progress had been made in developing workable gas dynamics computational models,[1-17] at the time this project was initiated, there was no available model which adequately predicted the mixed supersonic and subsonic flow regions present in the central core of the exhaust plumes. This report describes the results of a research program which was initiated to develop a finite difference computer code which would include the calculation of imbedded shock waves and Mach discs along with finite-rate chemistry and turbulent mixing phenomena.

During the course of this project, a gas dynamics computer code developed by Dr. Brian Spalding for the Rocket Propulsion Establishment in the United Kingdom was made available to the U.S. Army MIRADCOM as part of the TTCP program with Great Britain. This program was modified by the Principal Investigator to operate on the MIRADCOM computer. The results of a few limited calculations are described in this report.

The MIRADCOM Radiation Model (developed by H. Tracy Jackson, Jr. and extended by K. E. Harwell) was refined as part of the research effort described in this report. Results obtained using this model will be described and a comparison made between the theoretical predictions and experimental data.

2. THEORETICAL FOUNDATIONS

2.1 Gas Dynamics Theoretical Models

A complete review of the theory for turbulent exhaust plumes is beyond the scope of this report. A description of the flow theories is given by Jackson, Poslajko, and Harwell in Ref. 18,

The appropriate equations of motion for the gas are the conservation equations:

$$\text{Conservation of Mass: } \frac{\partial \rho}{\partial t} + \nabla \cdot \rho \vec{V} = 0$$

Conservation of Momentum:

$$\rho \left[\frac{\partial \vec{V}}{\partial t} + (\vec{V} \cdot \nabla) \vec{V} \right] = - \nabla p + \nabla \cdot \underline{\underline{\tau}}$$

Conservation of Energy:

$$\rho \frac{DE}{Dt} = \rho \frac{D \left[e + \frac{\vec{V} \cdot \vec{V}}{2} \right]}{Dt} = - \nabla \cdot (p \vec{V}) + \nabla \cdot \underline{\underline{\tau}} \cdot \vec{V} - \nabla \cdot \vec{q}$$

After supplementing these conservation equations with an Equation of State

$$e = e(T, p)$$

and constitutive relations relating the shearing stress τ to strain of the fluid element and ultimately to the velocity gradients in the flow, a complete system of flow equations is obtained. For example, the laminar shear stress can be expressed as

$$\tau_{ij} = \lambda \nabla \cdot \vec{V} \delta_{ij} + 2\mu \left(\frac{\partial V_j}{\partial X_i} + \frac{\partial V_i}{\partial X_j} \right)$$

The theoretical approach to the development of the turbulent shear stress model is not described here, but the model employed was the eddy diffusivity model proposed by Donaldson and Gray [19.] for the compressible free mixing of a primary jet with quiescent air:

$$\tau_t = -\bar{\rho} \bar{u}'\bar{v}' = K \bar{\rho} L U_s \frac{\partial \bar{u}}{\partial r}$$

where K is the mixing rate factor, L is a typical local scale length, $\bar{\rho}$ is the local density, and U_s is a typical velocity difference. These parameters are defined in Refs. [12 and 19].

It had been intended to incorporate the Turbulent Kinetic Energy (TKE) turbulence model approach into the computer model. However, due to problems with the basic computer code, this was not accomplished. In this model [see Ref.11], an additional conservation equation for the turbulent kinetic energy is added to the previous set of conservation equations. This equation for an axisymmetric flow becomes

$$\rho \bar{u} \frac{\partial \bar{k}}{\partial X} + \rho \bar{v} \frac{\partial \bar{k}}{\partial r} = \left[\frac{\rho}{P_{r_k}} \epsilon_t \frac{\partial \bar{k}}{\partial r} \right] + \rho \epsilon_t \left(\frac{\partial \bar{u}}{\partial r} \right)^2 - a_2 \frac{\bar{k}^{3/2}}{\ell_k}$$

where

$$\bar{k} = \frac{1}{2} \left[(\bar{u}')^2 + (\bar{v}')^2 \right]$$

$$\epsilon_t = \tau_t / \rho \frac{\partial \bar{u}}{\partial r}$$

$$\tau_t = a_1 \bar{\rho} \bar{k}$$

and

$$P_{r_k} = \bar{\rho} C_p \frac{\epsilon_t}{\bar{k}} . \quad -4-$$

a_1 and a_2 are Universal empirical relationships and λ_k is a length scale for the TKE. The reader is referred to Harsha [20] for a detailed formulation of this model. Suffice it to state that the addition of the TKE equation in the matrix formulation of the computer model is straightforward and can be accomplished without major difficulty.

The chemical reaction relations employed in the gas dynamics model was identical to those employed in the Low Altitude Plume Program (LAPP) developed for the U.S. Air Force Rocket Propulsion Laboratory [5]. The reader is referred to the LAPP report [5] for a detailed description of the formalism.

2.1.1 Gas Dynamics Finite Difference Computer Model

A finite difference numerical scheme was used to approximate the gas dynamics conservation equations described in the previous section.

A method developed by Daniel Matuska of the USAF Weapons Laboratory for computing hydrodynamic properties of fluids was modified to include the viscous stress tensor.

The finite differencing technique used was the standard one in which the derivative of the function F is defined as $[F(X + H) - F(X)]/H$ where H is finite.

The degree of success in using the finite difference technique to solve the differential equations depends on the "stability," "convergence" and "compatibility" of the method. The method is stable if any bounded starting procedure yields a uniformly bounded solution of the difference equation as H goes to zero. A method is convergent if the values of $Y(N)$, the solution of the finite difference expression approximating the difference equation, tend to the values of $Y(T)$ of the exact solution. The method is compatible

if the truncation error goes to zero as H goes to zero. The difference between the true solution of the differential equation and the finite difference equation method is influenced by the step size H and may require very small values of H to be stable and convergent.

One of the major advantages of the finite difference approach is that the problem is reduced to an initial value problem with well defined boundary conditions which results in a unique time-dependent solution. The major advantage in terms of the present problem of interest, is that the mixed subsonic/transonic/supersonic flow problem can be solved without using a patching solution.

The main disadvantages of the approach are that the complete boundary must be defined as a function of time, that small steps may be required which results in a fine computational mesh which in turn requires a large computer memory if realistic resolution is desired, and that computational times may become prohibitive.

The calculation procedure used in the present program is to treat the problem in two phases. In the first phase, the "Lagrangian terms" (coordinate system fixed to the fluid element) are integrated. In the second phase, the convective terms are integrated (mass momentum and energy is transported across boundaries). The details of the calculation procedure are described in Ref. [21].

2.2 Infrared Radiation Theoretical Model

A numerical infrared radiation band model was developed to predict the infrared spectral and spatial radiation intensity distributions in exhaust plumes. The model essentially extends the models of Jackson [22,23], which were developed for CO₂ radiation from aircraft plumes, to include water vapor radiation and to include the capability of treating higher temperatures than the Jackson code.

The statistical band model and spectral parameters used in the code were developed for application to rocket plume studies [24]. The transmissivity τ_R, in the model becomes

$$\tau_R = \exp [-P/(1 + P^2/4Q)^{1/2}]$$

where $P = \int k du$

$$Q = \int \frac{k\gamma}{d} du$$

$k(\omega, T)$ = absorption coefficient

u = optical depth

γ = average line width, and

$d(\omega, T)$ = average line spacing

The absorption coefficients k and reciprocal line spacings, $1/d$, are tabulated as a function of frequency, ω and temperature, T in Ref. [24].

The average line width was computed using the relation

$$\gamma = [120.1 P_{H_2O}/T] + [.09(P_{H_2O} + P_{N_2}) + .04 P_{O_2} + .12 P_{CO_2}] [273/T]^{1/2}$$

2.2.1 Infrared Radiation Computer Model

The infrared radiation computational model combined a water vapor emissivity formalism developed by Reardon [25] with the plume integration

scheme developed by Jackson [23] for carbon dioxide emission from aircraft plumes.

The computation of radiance from a jet exhaust plume was accomplished by dividing the plume into sufficiently small incremental blocks so that average conditions could be assigned and the blocks treated as homogeneous. The approach used in the geometric division of the plume is described in detail by Jackson [23]. The plume is subdivided into slabs parallel to the line of sight. Each slab is then cut into horizontal strips which lie along the line-of-sight vector.

The radiation from each of these strips is computed by integration using the modified Curtis-Godson approximation [24].

Integration is accomplished over the length of the strip to obtain the parameters P and Q as shown earlier,

$$P = \int kdu$$

$$Q = \int \frac{kY}{d} du$$

The transmittance $\tau_R = \exp [-P/(1 + P^2/4Q)^{1/2}]$ is then computed.

Radiance is then calculated by integrating over the strip emissivity. For a homogeneous block, Radiance = (Blackbody) (emissivity) so that the inhomogeneous strip radiance is

$$\text{STRIP RADIANCE} = \int (\text{Blackbody}) d\varepsilon$$

Since $d\varepsilon = -d\tau_R$ the computational algorithm

$$\text{STRIP RADIANCE} = - \sum_i (\text{Blackbody})_i \Delta \tau_{R,i}.$$

The algorithms used for calculating the radiation from inhomogeneous strips were tested by applying them to homogeneous strips and comparing with direct radiation calculations. The differences in calculated values were less than

one percent.

A listing of the infrared computer program is given in Appendix I.

A list of the input requirements is given in Appendix II. Appendix III includes a listing of a sample case with output.

3. DISCUSSION OF RESULTS

This section contains a discussion of the results obtained as part of the contract effort. The procedure used to validate the infrared radiation computer code is described in Section 3.1. Section 3.2 presents the results of the infrared radiation code predictions for a turbojet flight test for which there are experimental data. The prediction of the infrared radiation emitted from a small kerosene/gaseous oxygen rocket engine exhaust plume is presented in Section 3.3. A discussion of the progress in developing the Gas Dynamics Finite Difference Computer Model is described in Section 3.4. Some preliminary results obtained using the REP3 Computer Code are described in Section 3.5.

3.1 Validation of the UTSI Infrared Radiation Code

In order to validate the Infrared Radiation Code, the spectral radiance of water vapor and carbon dioxide mixtures experimental data are available [Refs. 24, 27-29]. Sukanek and Davis [26] performed a similar assessment of the NASA band model formulation developed by Reardon [25]. It was decided that several selected cases used by Sukanek and Davis would be used for the evaluation since this would provide a comparison of the present model calculations with the experimental data and with the predictions of the NASA/Reardon Radiation Code. The cases selected for the comparison are given in Table 3.1.

The spectral radiance of 60-cm thick slabs of hot H₂O vapor and CO₂ were computed for the 2.7- μ and 4.3- μ wavelength bands. Figures 3.1-3.8 present the results of the computations.

The comparison between the measured and predicted values of the spectral radiance in the 2.7- μ band for an isothermal slab of water vapor is presented

TABLE 3.1 CASES CONSIDERED FOR COMPARISON

Subject	Fig. No.	Particle Pressure (mmHg)					Temperature Profile	Length (m)
		H ₂ O	CO ₂	N ₂	CO	H ₂		
Isothermal Hot H ₂ O	1	38.0		38.7			1215	0.6
Isothermal Hot CO ₂	2		38.0	38.0			1202	0.6
Isothermal Hot CO ₂	3		7.58	68.43			1200	0.6
Isothermal Mixture	4	35.26	7.58	27.09	1.44	4.64	1200	0.6
Isothermal Mixture	5	35.26	7.58	27.06	1.47	4.62	1200	0.6
Nonisothermal H ₂ O	6	670					382/537/723/953/1128/1160/990/751/558/389	0.6
Nonisothermal CO ₂	7		760				386/528/719/953/1130/1160/975/737/541/387	0.6
Nonisothermal Mixture	8	57	28	675			378/537/723/958/1127/1158/990/752/555/383	0.6

NOTE: The temperature of the gas was either homogenous, in which case one temperature is reported, or was inhomogeneous, in which case the temperatures are given at the center of each 6 cm wide cell.

in Fig. 3.1. At the center of the spectral band the computed values are slightly lower than the measured ones. The NASA code predicts slightly higher values than the UTSI code, but in general the agreement is fairly good between the calculated and measured values. The calculated values are within 20% of the measured ones.

The measured and predicted values of spectral radiance in the $2.7-\mu$ and $4.3-\mu$ wavelength bands for an isothermal slab of CO_2 are presented in Figs. 3.2 and 3.3. In contrast to the H_2O predictions, the predictions of the spectral radiance for CO_2 do not agree very well with the measurements. Errors of 30 to 50 percent are observed. Similar disagreement was found by Lindquist and others [27-29] who attributed to the probable cause of the disagreement to the lack of knowledge of the CO_2 broadening parameters at low pressures. Ludwig and others [24] show better agreement between the predictions and measurements at higher pressures ($P_{\text{CO}_2} = 1-2$ atmospheres).

Figures 3.4 and 3.5 present the results for the spectral radiance of an isothermal mixture of CO_2 and H_2O in the 4.3 and $2.7-\mu$ bands, respectively. The calculated values are higher than the measurements in the center of the spectral bands. Both codes yield results which are in fair agreement with the measured values.

Figure 3.6 presents the results for a nonisothermal slab of water vapor. There is fairly good agreement between the predicted and measured values, but the present code appears to predict values which are below the measured and calculated values using the NASA code.

Figure 3.7 presents the comparison between the predicted and measured values of spectral radiance in the $2.7-\mu$ spectral band of a nonisothermal slab of CO_2 . The agreement between the predicted and measured values is quite good.

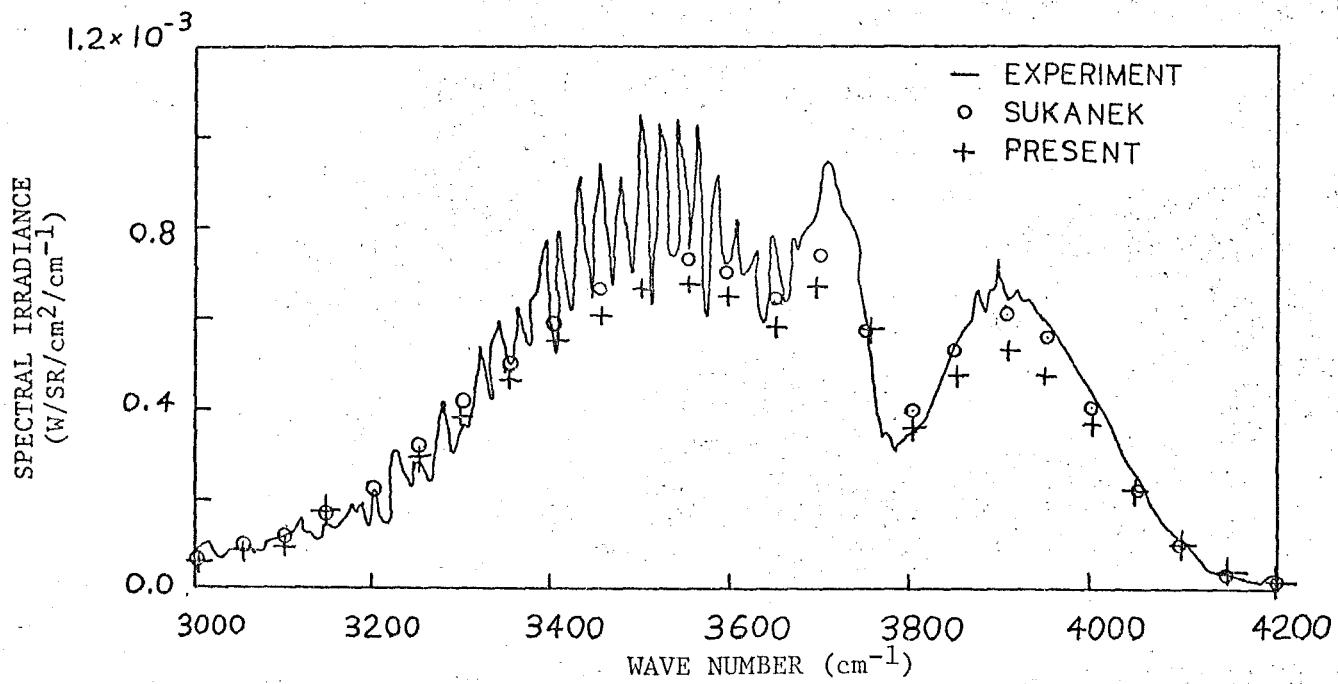


Fig. 3.1 SPECTRAL RADIANCE IN THE $2.7-\mu\text{m}$ BAND FOR AN ISOTHERMAL SLAB OF H_2O VAPOR.

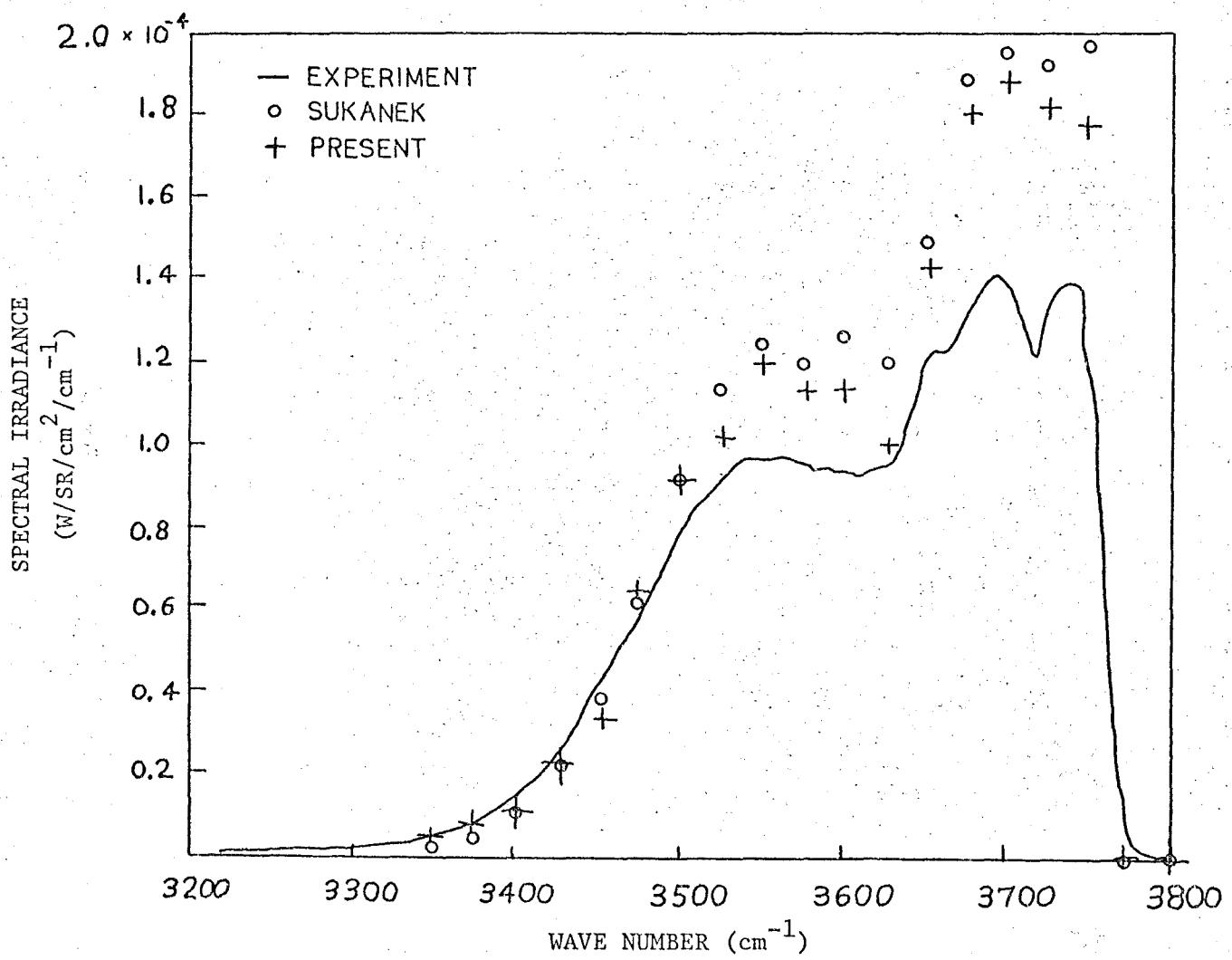


Fig. 3.2 SPECTRAL RADIANCE IN THE 2.7- μm BAND FOR AN ISOTHERMAL SLAB OF CO_2 .

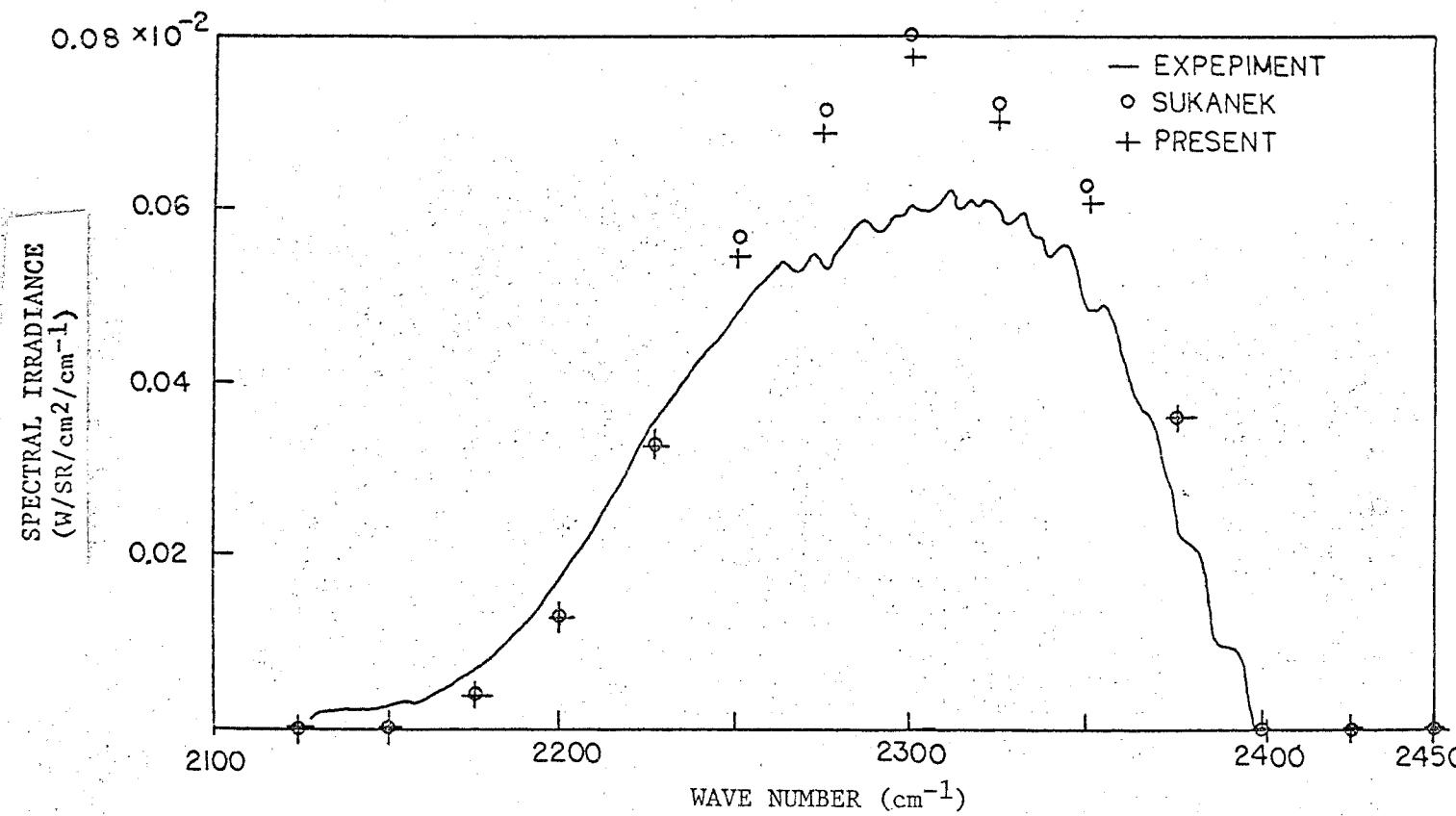


Fig. 3.3 SPECTRAL RADIANCE IN THE 4.3- μ m BAND FOR AN ISOTHERMAL SLAB OF CO₂.

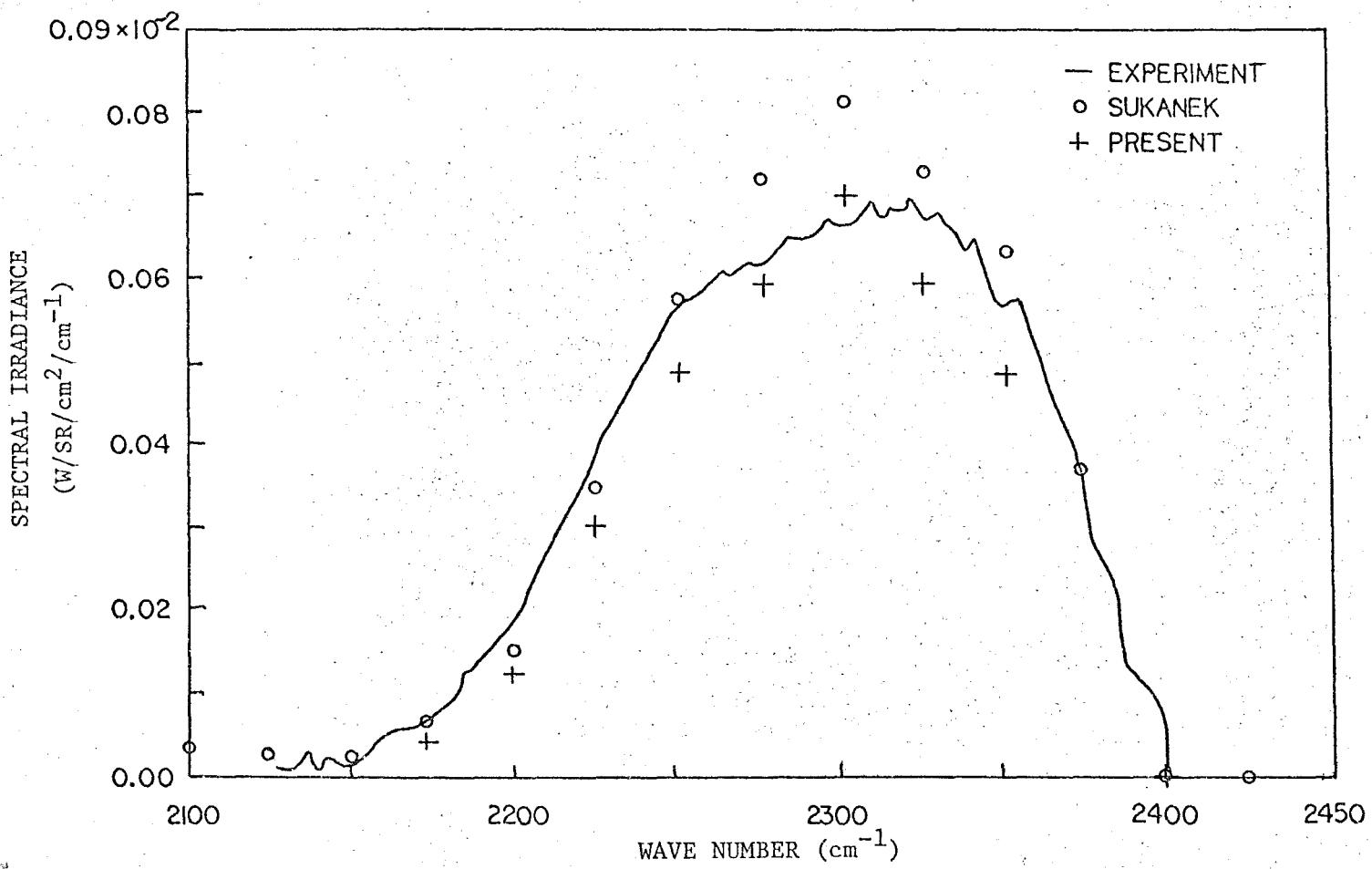


Fig. 3.4 SPECTRAL RADIANCE IN THE 4.3- μm BAND FOR AN ISOTHERMAL MIXTURE OF CO_2 and H_2O .

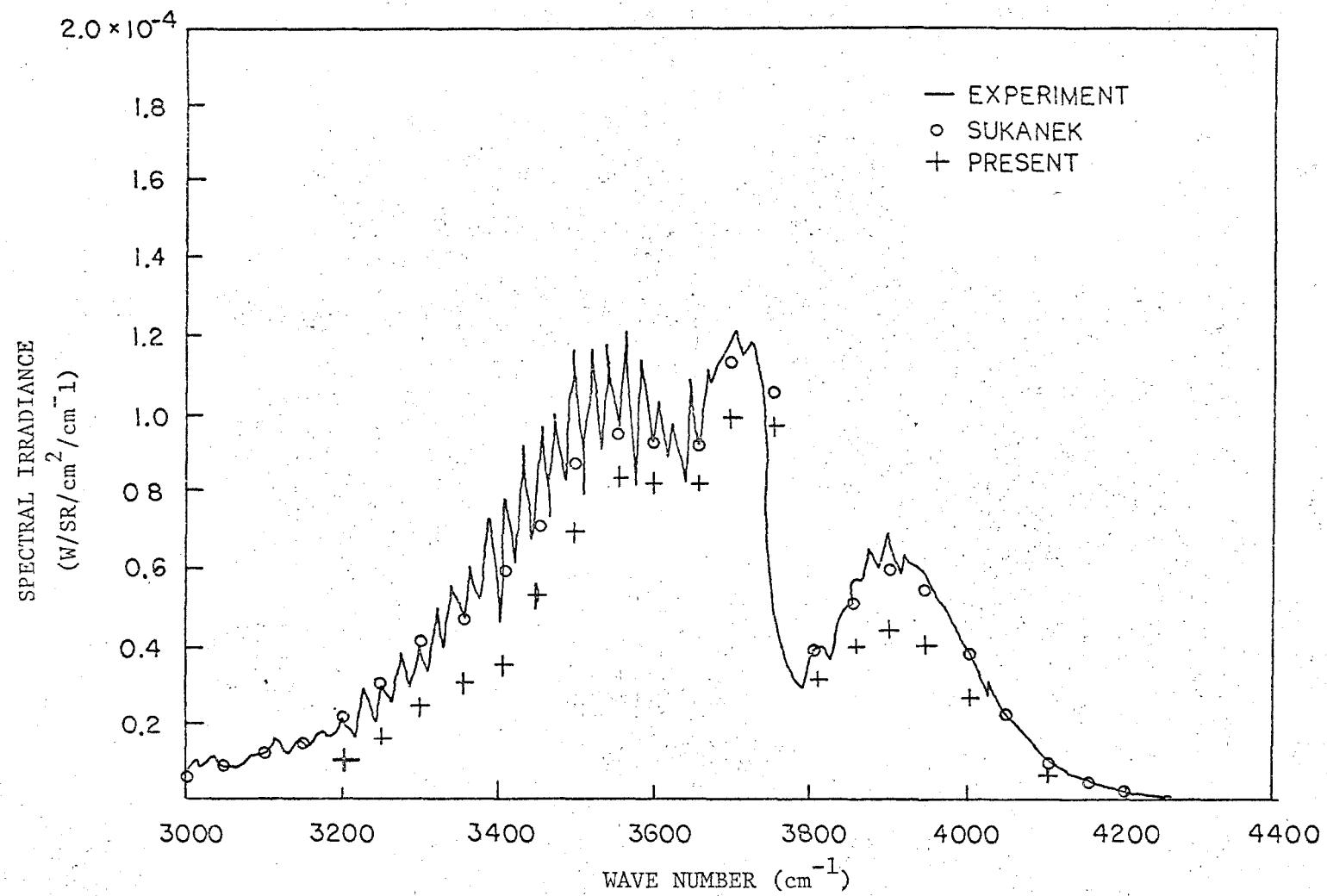


Fig. 3.5 SPECTRAL RADIANCE IN THE 2.7- μm BAND FOR AN ISOTHERMAL MIXTURE OF CO₂ AND H₂O.

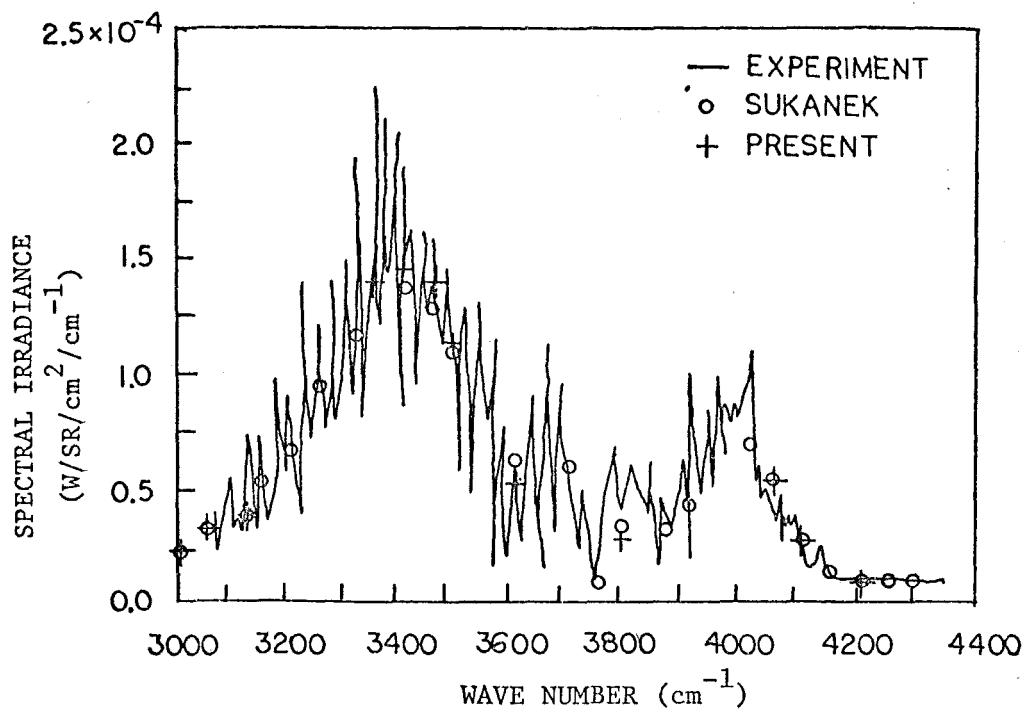


Fig. 3.6 SPECTRAL RADIANCE IN THE 2.7- μm BAND FOR A NONISOTHERMAL SLAB OF H_2O VAPOR.

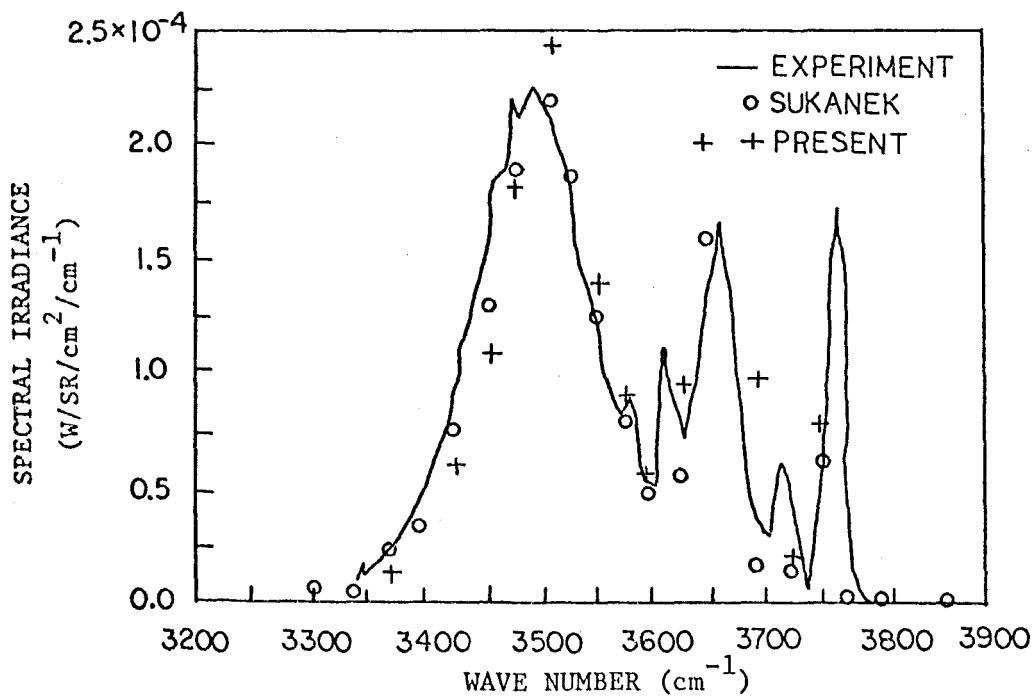


Fig. 3.7 SPECTRAL RADIANCE IN THE 2.7- μm BAND FOR A NONISOTHERMAL SLAB OF CO_2 .

Figure 3.8 presents the results for a nonisothermal mixture of CO_2 and H_2O . There is fairly good agreement over the entire $2.7-\mu$ wavelength band.

In summary, the UTSI radiation code appears to yield results which agree fairly well with the NASA code even though there are some differences. With the exception of the low pressure CO_2 cases, the predicted values are within $\pm 20\%$ of the measured values.

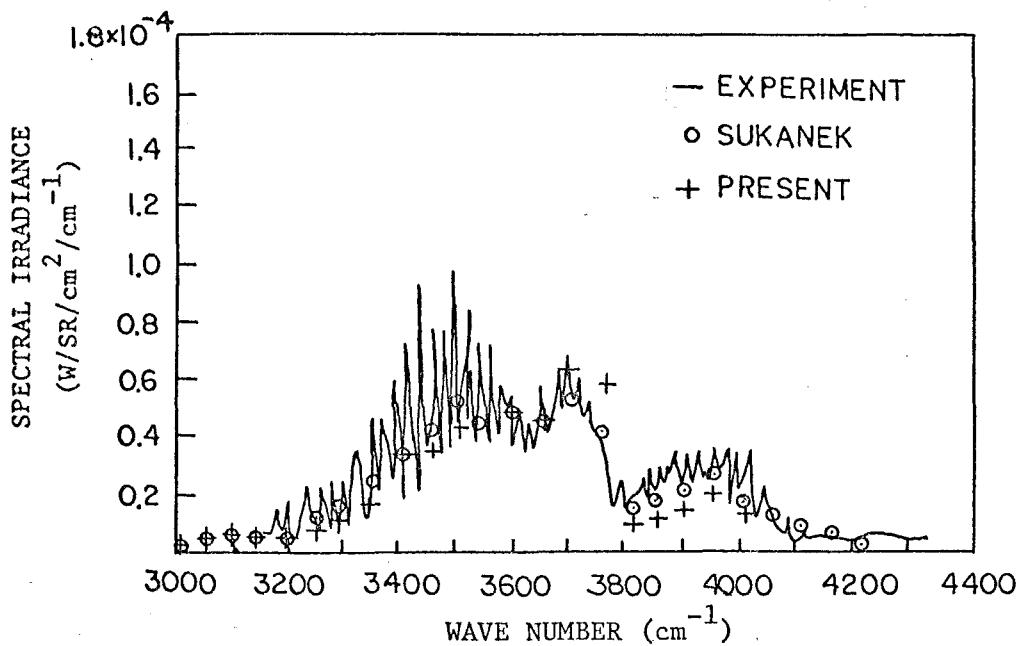


Fig. 3.8 SPECTRAL RADIANCE IN THE $2.7-\mu\text{m}$ BAND FOR A NONISOTHERMAL MIXTURE OF CO_2 AND H_2O .

3.2 Infrared Radiation Predictions for Turbojet Flight Test

The UTSI Infrared Radiation Code was utilized to calculate the infrared radiation emitted in the wavelength interval from 4 to 5 microns from a turbojet aircraft exhaust for which experimental data are available.

H. Tracy Jackson, Jr. at MIRADCOM provided the data given in Table 3.2.

The Low Altitude Plume Program (LAPP) was exercised for the input data given in Table 3.2. The first portion of the LAPP code output for this case, labeled "FLT1-PASS 7," is included as Figs. 3.9-3.11. The output data from the LAPP code was then used as input data for the Infrared Radiation Code. Refer to Fig. 3.12 for a typical set of input data.

The results of the computations are presented in Figs. 3.13-3.15. The CO_2 contribution to the radiation per length of plume (Watts/SR/cm) is given in Fig. 3.13 while the H_2O contribution is given in Fig. 3.14. The H_2O radiation is very small compared to the CO_2 radiation.

The spatial distribution of the radiation is given in Fig. 3.15. Also, shown on the figure is the centerline temperature distribution. The total radiation emitted from the plume in the 4.0 to 5.0 micron wavelength interval was 165.7 Watts/SR. This radiation is at the plume and is uncorrected for detector response characteristics or for atmospheric attenuation between the plume and the detector.

It had been anticipated that the Jackson MIRADCOM Radiation Code would be used to calculate the radiation for the same test flight so that a comparison could be made between the two codes. Unfortunately, the results from the MIRADCOM Code did not become available prior to the completion of the project.

TABLE 3.2

TURBOJET FLIGHT TEST DATA

Flight 10, Pass 7

TPT = 790° K

Vel = 512 Kts. = 263 m/sec.

Asp. α = 19°

O/F = 62

Range ~ 2212 meters

Too = 526°R

 ρ_{∞} = 1955 lb/ft²

C* = 1124 ft/sec.

Air Flow 182 lb/sec., 99% RPM

V jet = 1801.6 ft/sec.

AEROCHEM RESEARCH LABORATORIES PRINCETON N.J.
A FAST COMPUTER PROGRAM FOR NONEQUILIBRIUM ROCKET PLUME PREDICTIONS
THIS VERSION INCLUDES SHOOT AND MODEL 6 MODIFICATIONS

F_LFI - PASS 7

PRESSURE(CONSTANT) = 1.000000E 00 ATMOSPHERES

NOZZLE RADIUS= 8.9999998E-01 FEET

DEVIS NUMBER(CONSTANT)= 1.000000E 00 PRANDTL NUMBER(CONSTANT)= 6.999999E-01

X INITIAL(FEET)= 0.0

X FINAL(FEET)= 3.000000E 01

DT/F INCREMENT= 1.000000E 00

MINIMUM STEP SIZE= 9.9999994E-11

DRILLIUSON/GRAY VISCOSITY MODEL

	JET	EDGE
TEMPERATURE(DEG. KELVIN)	7.900000E 02	2.920000E 02
VELOCITY (FEET/SECOND)	1.800000E 03	1.000000E 02
MOLE FRACTION H	2.8829817E-07	2.8849433E-07
MOLE FRACTION H2	2.8829817E-07	2.8849433E-07
MOLE FRACTION H2O	3.3150621E-02	2.8849433E-07
MOLE FRACTION CO	2.8829817E-07	2.8849433E-07
MOLE FRACTION CO2	3.2630604E-02	3.2879249E-04
MOLE FRACTION O	2.8829817E-07	2.8649433E-07
MOLE FRACTION OH	2.8829817E-07	2.8849433E-07
MOLE FRACTION O2	1.5670262E-01	2.0846844E-01
MOLE FRACTION N2	7.7751452E-01	7.9120100E-01

	REACTIONS BEING CONSIDERED	KR=A*EXP(B/RT)/T**N	A	N	B	(MOLECULE-ML-SEC UN)
1	O + O + M = O2 + M	1.000E-29	1.0	0.0		
2	H + H + M = H2 + M	1.000E-29	1.0	0.0		
3	O + H + M = OH + M	1.000E-29	1.0	0.0		
4	H + O2 + M = H2O + M	1.000E-28	1.0	0.0		
5	CO + O + M = CO2 + M	5.000E-29	1.0	-4000.0		
6	OH + O2 = H2O + O	1.000E-11	0.0	-1000.0		
7	OH + H2 = H2O + H	4.000E-11	0.0	-5500.0		
8	O + H2 = OH + H	3.000E-11	0.0	-8200.0		
9	H + O2 = OH + O	3.000E-10	0.0	-16500.0		
10	CO + O2 = CO2 + O	5.000E-13	0.0	-600.0		

Fig. 3.9 Initial Conditions for LAPP Computer Program

0.0

FEET

PUT1-PASS7

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X/R DELTA X FEET PRESS(ARM)
 0.0 8.999491E-02 1.000000E 00

HALF RADIUS/R 1.13722E-00	INNER MIXING ZONE RADIUS/R 1.013722E-00	MACH NUMBER AT HALF RADIUS 6.269119E-01	MIXING RATE COEFFICIENT 4.209092E-02
---------------------------------	---	---	--

Y/R	VELOCITY FEET/SEC	TEMPERATURE K	DENSITY GM/CC	MACH NO.	ENTHALPY CAL/GM	VISCOSITY LB/FT/SEC	ELECTRON DENSITY(1/ML)	PSI
0.0	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	0.0	
3.1900	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	1.121742E-01	
3.2000	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	2.243484E-01	
3.3000	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	3.365226E-01	
3.4000	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	4.486969E-01	
3.5000	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	5.608711E-01	
3.6000	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	6.730453E-01	
3.7100	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	7.852195E-01	
3.8000	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	8.973937E-01	
3.9000	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	1.009567E 00	
4.0000	1.800000E-03	7.900000E-02	4.448628E-04	9.920356E-01	-4.688548E-01	1.104238E-01	1.121741E 00	
3.2745	1.000000E-02	2.920000E-02	1.204586E-03	8.888811E-02	-2.495791E-00	2.989527E-01	1.233916E 00	
3.8338	1.000000E-02	2.920000E-02	1.204386E-03	8.888811E-02	-2.495791E-00	2.989527E-01	1.345090E 00	

Fig. 3.10 Sample Output of LAPP Computer Program

CODE EXTRACTS

	H ₂ O	CO	C ₂ O ₂	D	OH	PI
42	2.68298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	1
43	2.83298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2
44	2.83298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2
45	2.83298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	3
46	-	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	-
47	-	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	-
48	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	4
49	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	5
50	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	6
51	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	7
52	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	8
53	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	9
54	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	10
55	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	11
56	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	12
57	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	2.88298E-07	13

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	H2	H2J	CO	CO2	O	DH	PT
-5.29371E-04	5.31415E-10	-1.39146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	1
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	2
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	3
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	4
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	5
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	6
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	7
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	8
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	9
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	10
-5.29371E-04	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	-6.40717E-08	6.87579E-08	11

Fig. 3.11 Sample Output of LAPP Computer Program

= 0.0 FEET

FUJI-PASS7

PAGE 1

MOLE FRACTIONS

<i>t/k</i>	0.2	H ₂	B	SRMO	MED	2	PI
0.0	1.55703E-01	7.77515E-01	0.0	0.0	0.0	0.0	1
0.10000	1.36703E-01	7.77515E-01	0.0	0.0	0.0	0.0	2
0.20000	1.22703E-01	7.77515E-01	0.0	0.0	0.0	0.0	3
0.30000	1.103E-01	7.77515E-01	0.0	0.0	0.0	0.0	4
0.40000	1.26703E-01	7.77515E-01	0.0	0.0	0.0	0.0	5
0.50000	1.26703E-01	7.77515E-01	0.0	0.0	0.0	0.0	6
0.55700	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	7
0.70000	1.359703E-01	7.77515E-01	0.0	0.0	0.0	0.0	8
0.80000	1.36703E-01	7.77515E-01	0.0	0.0	0.0	0.0	9
0.90000	1.26703E-01	7.77515E-01	0.0	0.0	0.0	0.0	10
1.00000	1.26703E-01	7.77515E-01	0.0	0.0	0.0	0.0	11
1.27445	2.042594E-01	7.91251E-01	0.0	0.0	0.0	0.0	12
1.63353	2.08486E-01	7.91251E-01	0.0	0.0	0.0	0.0	13

NET RATE OF PRODUCTION (W-DOT/KHO\$U)

<i>y₂</i>	<i>y₂</i>	<i>B</i>	SKMO	MED	2	PI
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	1
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	2
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	3
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	4
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	5
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	6
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	7
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	8
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	9
-6.34503E-08	0.0	0.0	0.0	0.0	0.0	10
LESS THAN 1E-08	0.0	0.0	0.0	0.0	0.0	11

Fig. 3.11 (cont.)

PRESSURE	EXIT RADIUS	T(K)	P <small>N₂</small>	P <small>CO</small>	P <small>CO₂</small>	P <small>H₂O</small>	P <small>O₂</small>
.1000E 010	.2720E 02	120					
0.0 = X POSITION(cm)	R(cm)						
0.0	0.7900E 030	1.000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.2720E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.5440E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.8160E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1088E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1360E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1632E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1904E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.2176E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.2448E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.2720E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.3400E	020.2920E	030.1000E-190	1.000E-190	3.288E-030	2.885E-060	1.000E-19	
0.382E	04 = X(cm)						
0.0	0.7900E 030	1.000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.2720E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.5440E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.8160E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1088E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1360E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1632E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1904E	020.7899E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.2176E	020.7876E	030.1000E-190	1.000E-190	3.245E-010	3.296E-010	1.000E-19	
0.2448E	020.7640E	030.1000E-190	1.000E-190	3.061E-010	3.108E-010	1.000E-19	
0.2728E	020.6661E	030.1000E-190	1.000E-190	2.319E-010	2.346E-010	1.000E-19	
0.3123E	020.4492E	030.1000E-190	1.000E-190	9.216E-020	9.121E-020	1.000E-19	
0.3968E	020.3143E	030.1000E-190	1.000E-190	1.604E-020	1.309E-020	1.000E-19	
0.5252E	020.2935E	030.1000E-190	1.000E-190	4.128E-030	8.655E-040	1.000E-19	
0.6351E	020.2920E	030.1000E-190	1.000E-190	3.288E-030	2.885E-060	1.000E-19	
0.720E	02						
0.0	0.7900E 030	1.000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.2720E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.5440E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.8160E	010.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1088E	020.7900E	030.1000E-190	1.000E-190	3.263E-010	3.315E-010	1.000E-19	
0.1360E	020.7899E	030.1000E-190	1.000E-190	3.263E-010	3.314E-010	1.000E-19	
0.1632E	020.7892E	030.1000E-190	1.000E-190	3.257E-010	3.309E-010	1.000E-19	
0.1903E	020.7848E	030.1000E-190	1.000E-190	3.222E-010	3.273E-010	1.000E-19	
0.2175E	020.7667E	030.1000E-190	1.000E-190	3.078E-010	3.125E-010	1.000E-19	
0.2449E	020.7179E	030.1000E-190	1.000E-190	2.694E-010	2.731E-010	1.000E-19	
0.2744E	020.6259E	030.1000E-190	1.000E-190	2.022E-010	2.041E-010	1.000E-19	

Fig. 3.12 Typical Set of Input Data for Radiation Code

FL1.PASS7

X(cm)

W/SR/CM

STATION RADIATION (W/SR)

0.0	0.8030024E 00	0.1062637E 01
38.19	0.6650596E 00	0.2803171E 02
71.97	0.5855711E 00	0.2112463E 02
102.27	0.5204300E 00	0.1675532E 02
135.10	0.4548181E 00	0.1600888E 02
165.70	0.3968118E 00	0.1302993E 02
195.02	0.3400359E 00	0.1190510E 02
228.89	0.2897733E 00	0.9721726E 01
259.24	0.2418252E 00	0.8068358E 01
298.38	0.1936032E 00	0.8521406E 01
329.39	0.1511664E 00	0.5345317E 01
359.80	0.1196913E 00	0.412519E 01
395.49	0.9243929E-01	0.3779300E 01
429.76	0.7282054E-01	0.2831890E 01
462.94	0.5788064E-01	0.2166581E 01
494.77	0.4560884E-01	0.1652521E 01
525.78	0.3838268E-01	0.1317710E 01
560.59	0.3459402E-01	0.1270374E 01
594.32	0.2598157E-01	0.1021555E 01
626.96	0.2198615E-01	0.7828305E 00
658.51	0.1888854E-01	0.6446391E 00
689.25	0.1645325E-01	0.5431341E 00
723.25	0.1424051E-01	0.5217939E 00
755.43	0.1251272E-01	0.4438905E 00
786.80	0.11049931E-01	0.3821360E 00
820.08	0.9954128E-02	0.3292734E 00
850.82	0.1471464E-01	0.3791106E 00
884.27	0.8120896E-02	0.3819935E 00
909.84	0.0	0.1038170E 00

* RADIATION VALUE WAS CALCULATED BASED ON THE SLOPE BETWEEN THE PREVIOUS TWO POINTS

TOTAL RADIATION EMITTED=0.1612E 03 WATTS/STER

CENTROID= 0.1921E 03C

Fig. 3.13 CO₂ Contribution to Infrared Radiation per Length of Plume

FL1.PASS7

STATION RADIATION (W/SR)

0.0	0.2072786E-01	0.1062637E 01
38.19	0.1714092E-01	0.7230808E 00
71.97	0.1537959E-01	0.5493103E 00
102.27	0.1386401E-01	0.4430514E 00
135.10	0.1225346E-01	0.4287230E 00
165.70	0.1080529E-01	0.3527986E 00
198.02	0.9366962E-02	0.3259194E 00
228.89	0.8083675E-02	0.2693681E 00
259.24	0.6870915E-02	0.2269739E 00
298.38	0.5615387E-02	0.2443608E 00
329.39	0.4495839E-02	0.1567802E 00
359.86	0.3648198E-02	0.1240556E 00
395.49	0.2844438E-02	0.1165630E 00
429.76	0.2336578E-02	0.8953859E-01
462.94	0.1902269E-02	0.7033050E-01
494.77	0.1586058E-02	0.5518754E-01
525.78	0.1313538E-02	0.4464533E-01
560.59	0.1114177E-02	0.4220155E-01
594.32	0.9213609E-03	0.3432758E-01
625.96	0.7910719E-03	0.2794681E-01
658.51	0.6881284E-03	0.2333586E-01
699.25	0.6061282E-03	0.1989019E-01
733.25	0.5306695E-03	0.1932555E-01
756.43	0.4710052E-03	0.1661982E-01
783.80	0.4216461E-03	0.1444663E-01
820.03	0.3811687E-03	0.1255593E-01
850.82	0.3897771E-03	0.1184792E-01
884.27	0.3154260E-03	0.1179666E-01
909.84	0.0	0.4032392E-02

* RADIATION VALUE WAS CALCULATED BASED ON THE SLOPE BETWEEN THE PREVIOUS TWO POINTS

TOTAL RADIATION EMITTED=0.4459E 01 WATTS/STER

CENTERED= 0.2050E 03CM

Fig. 3.14 H₂O Contribution to Infrared Radiation per Length of Plume

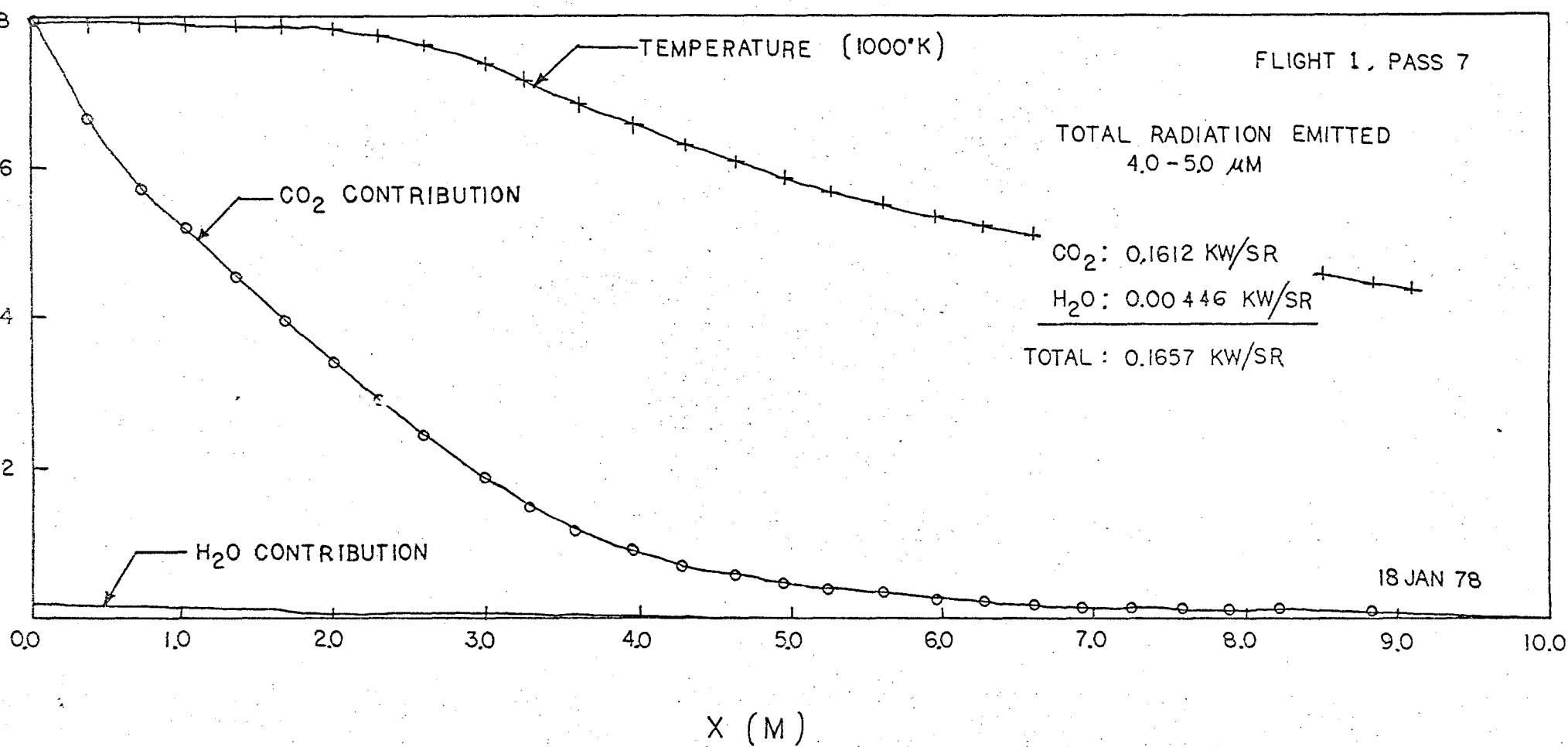


Fig. 3.15 SPATIAL DISTRIBUTION OF TEMPERATURE AND RADIATION IN THE TURBOJET EXHAUST.

3.3 Infrared Radiation Predictions for a Small Kerosene/Gaseous Oxygen Rocket Engine

Two gas dynamics and the UTSI infrared radiation computational models were used to predict the physical properties in hot rocket exhaust plumes. Experimental data were obtained in the exhaust plume of a small kerosene/gaseous oxygen rocket engine using a three-dimensional laser Doppler velocimeter system to measure the velocity distribution. The spatial distribution of radiated energy in the $4-5 \mu$ wavelength band was measured using an infrared radiometer. Detailed comparisons between the computational and experimental results are described in this section.

3.3.1 Description of Experimental Equipment

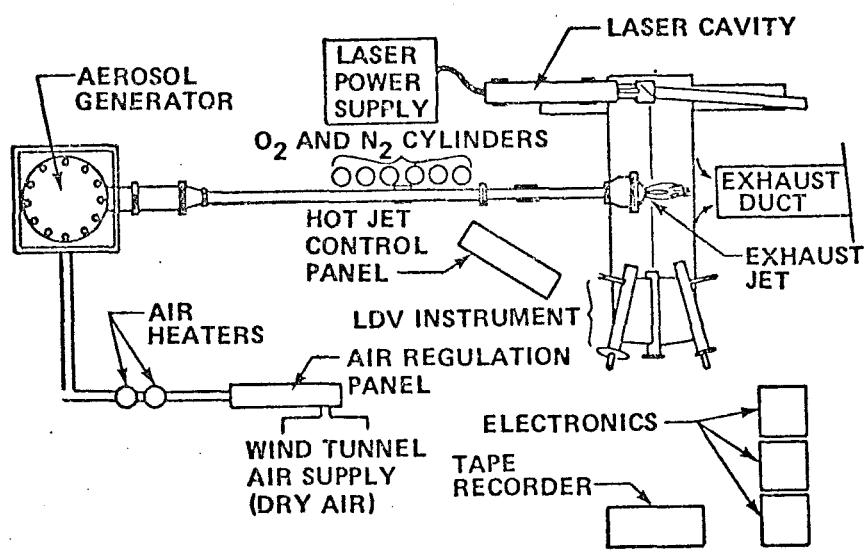
Small Laboratory Rocket Engine: The exhaust plumes studied in these measurements were produced by a small kerosene/gaseous oxygen rocket engine [30] operating in the NASA Laser Doppler Velocimeter Facility [31]. Engine design and performance parameters are listed in Table 3.3. Data were taken at a variety of oxidizer/fuel (O/F) ratios and at two chamber pressures. The experimental data used to compare with the theoretical calculations were obtained at a chamber pressure of 11.1 atm. and an O/F ratio of 2.25.

NASA Laser Doppler Velocimeter Facility: The NASA Marshall three-dimensional laser Doppler velocimeter [31] was used to measure the velocity distribution in the exhaust plume and mixing layer produced by the inner hot rocket exhaust jet mixing with an outer subsonic airflow. A schematic diagram of the LDV experimental arrangement showing the installation of the small rocket engine is shown in Fig. 3.16. As described in Ref. [18] the simultaneous measurement of the three Doppler velocity components yielded the three velocity components u, v , and w .

TABLE 3.3

ENGINE DESIGN AND PERFORMANCE PARAMETERS

Fuel	Kerosene
Oxidizer	Gaseous oxygen
Purge	Nitrogen 100 psig
Coolant	Water at 3 gal/min
Chamber pressure	50 and 148 psig
Fuel flow rate	
50 psig	1.38 g/sec
148 psig	3.80 g/sec
Oxygen flow rate	
50 psig	3.05 g/sec
148 psig	8.70 g/sec
Throat diameter	0.200 in.
Exit diameter	0.310 in.
Throat area	0.0314 in. ²
Exit area	0.0755 in. ²
Area ratio (A_E/A_t)	2.40
O/F ratio	2.25
Chamber diameter	0.780 in.
Chamber length	1.620 in.
L^*	25.20 in.
C^* at 148 psig	6045 fps



3.16 SCHEMATIC DIAGRAM OF LASER VELOCIMETER ARRANGEMENT

In order to simulate external airflow around the small rocket engine, an outer flow nozzle was used to produce an external flow velocity of 61 m/sec. A schematic diagram of the rocket engine and outer flow nozzle is presented in Fig. 3.17.

Infrared Radiation Diagnostic Equipment: The primary instrument used to obtain the radiation intensity data was an Electro-Optical Industries Model 470 Radiometer. The data were obtained using a fixed filter, which transmitted radiation in the 4-5 μ band.

A large metal shield with a 1-mm aperture was moved in the axial direction to obtain the variation of radiant intensity along the exhaust plume centerline. A second set of measurements was made using a 1-mm vertical slot in the large metal shield. The slot length was greater than the diameter of the exhaust plume. The experimental setup is shown in Fig. 3.18.

3.3.2 Typical Experimental Results

Velocity Distribution Data: Experimental data for mean flow velocities ($\bar{u}, \bar{v}, \bar{w}$), turbulent intensities [$(\bar{u'})^2, (\bar{v'})^2, (\bar{w'})^2$], and turbulent velocity correlations were obtained in the radial direction at various axial locations downstream from the nozzle exit plane. Data were obtained at X/D locations of 2.4, 4.8, 8.4, 11.3, and 14.2, where X is the distance measured from the nozzle exit plane and D is the diameter of the exit nozzle. Typical data obtained at X/D = 11.3 for the mean axial flow velocity are presented in Fig. 3.19. Other data are described in Ref. [18].

Infrared Radiation Data: The infrared radiation data have been described previously [30]. Typical data obtained at an O/F ratio of 2.25 and an external flow velocity of 210 fps are presented in Figs. 3.20 and 3.21 for the 1-mm aperture and slot, respectively. The centerline radiation intensity

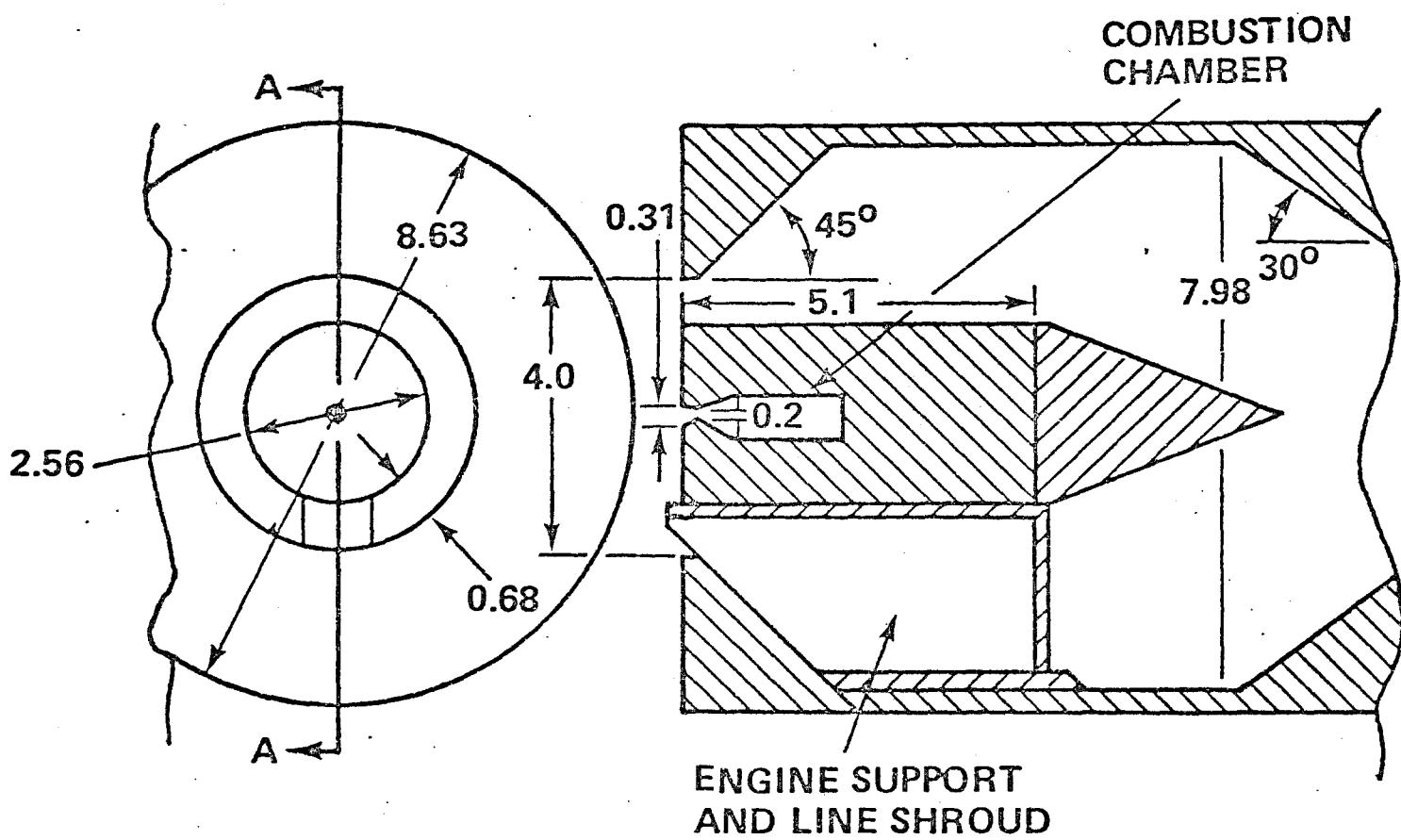


Fig. 3.17 Schematic of the exhaust plume and external flow apparatus.

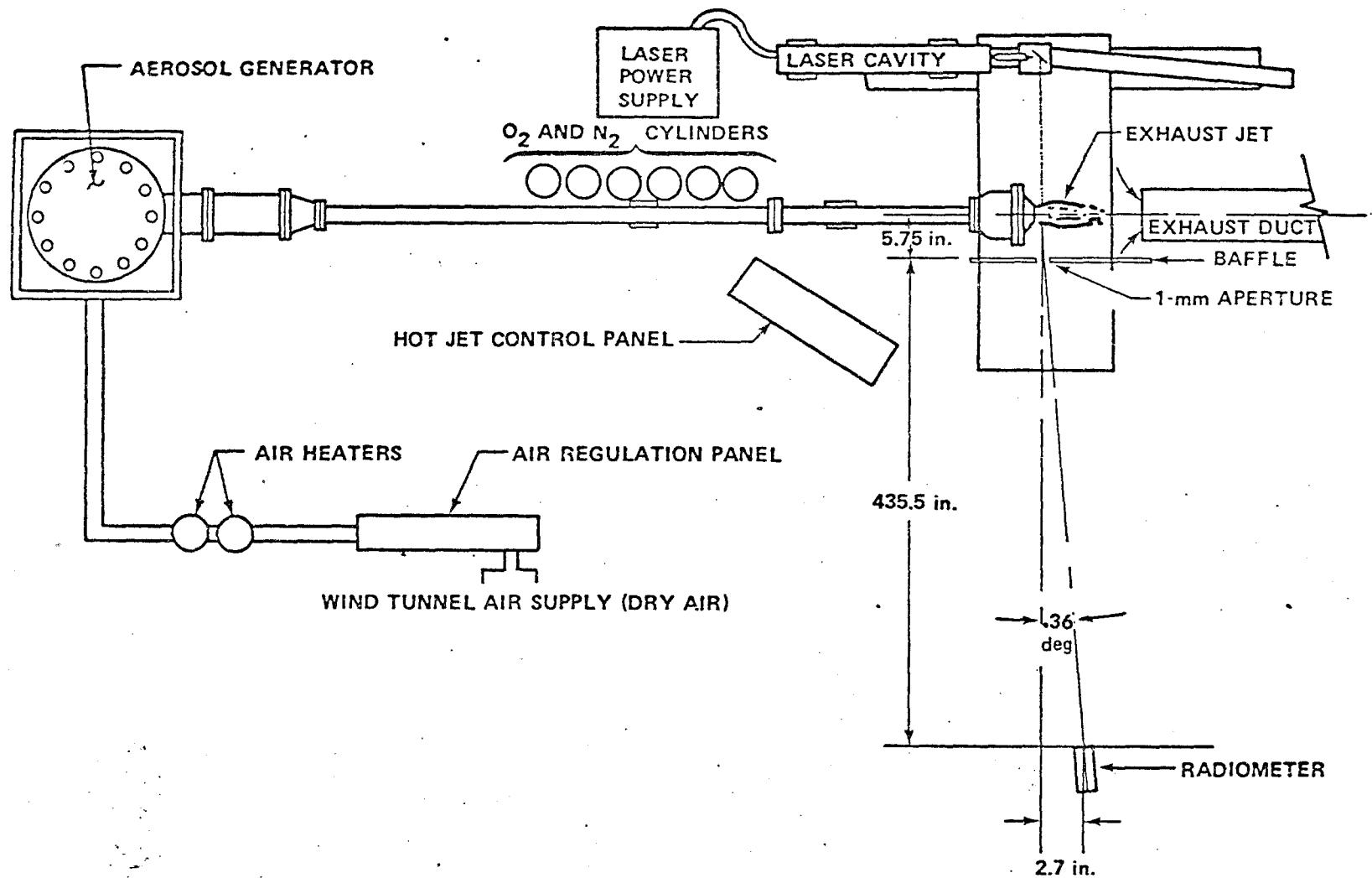


Fig. 3.18 Exhaust plume test arrangement in NASA LDV facility.

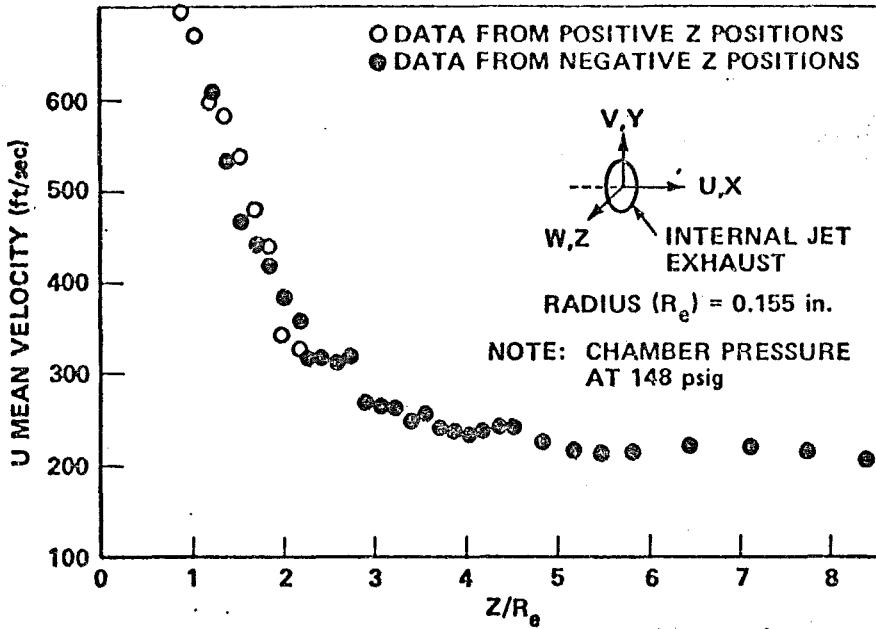


Fig. 3.19. Profile of u mean velocity component at $X/D = 11.3$

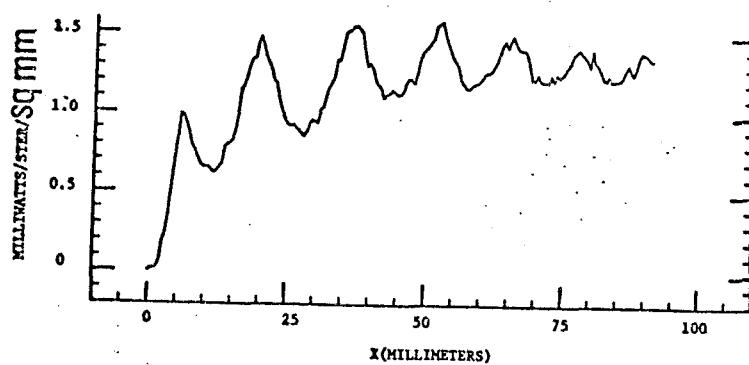


Fig. 3.20 Spatial distribution of radiation along axis of plume for aperture diameter = 1-mm; O/F = 2.25, velocity = 150 fps.

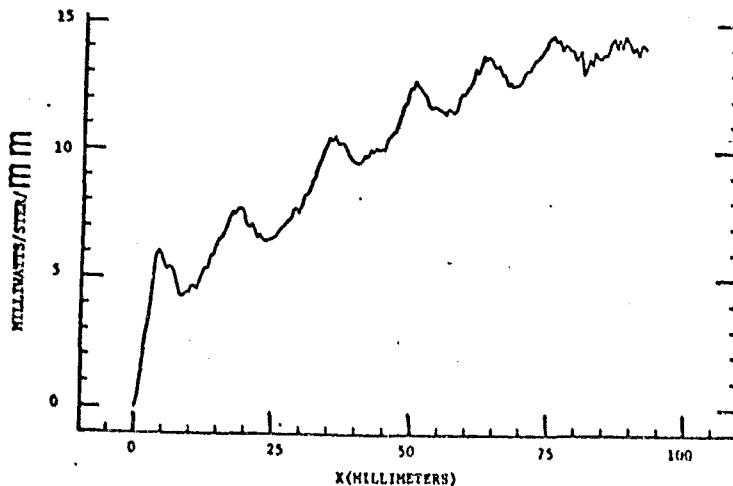


Fig. 3.21 Spatial distribution of radiation along axis of plume for 1-mm slot width; O/F = 2.25, velocity = 210 fps.

distribution in Fig. 3.20 clearly shows the presence and influence of shocks and/or Mach disks in the exhaust plume. The effect of afterburning in the mixing layers on the infrared signature is shown clearly in Fig. 3.21 where the shock-induced radiation peaks are superimposed on an axially increasing radiation intensity. The data presented in Figs. 3.20 and 3.21 were obtained in the near field of the exhaust plume (X/D less than 12) and indicate that the peak radiation intensity has not been reached. Typically, the axial distribution of radiation is similar to that shown in Fig. 3.22. Even though the data presented in Fig. 3.22 are for a lower mass flow and chamber pressure, the axial variation is similar to that obtained for the high-chamber-pressure cases. The data presented here are only samples selected to indicate the measured trends in preparation for a discussion of the comparison between theory and experiment.

3.3.3 Comparison of Calculated and Measured Results

The LAPP and TKE computer codes were exercised for the rocket performance parameters listed in Table 3.3. The NASA Lewis Chemical Equilibrium Computer Program was used to calculate the exit plane gas properties. Because of the short expansion nozzle length, the chemical species were assumed to be frozen at the throat values. These exit plane conditions then were used as initial conditions for the LAPP and TKE computer codes.

The results predicted by the LAPP and TKE computer codes are presented in Figs. 3.23-3.28. The variation of the centerline velocity with distance from the nozzle exit plane is presented in Fig. 3.23. The LAPP code predicts a constant-velocity core having a length approximately equal to $X/D = 5$, which is followed by an exponential velocity decay due to turbulent mixing predicted by the Donaldson-Gray viscosity model. The TKE (REP3) code predicts a sharp decrease in velocity at approximately an S/D value of one. There are also

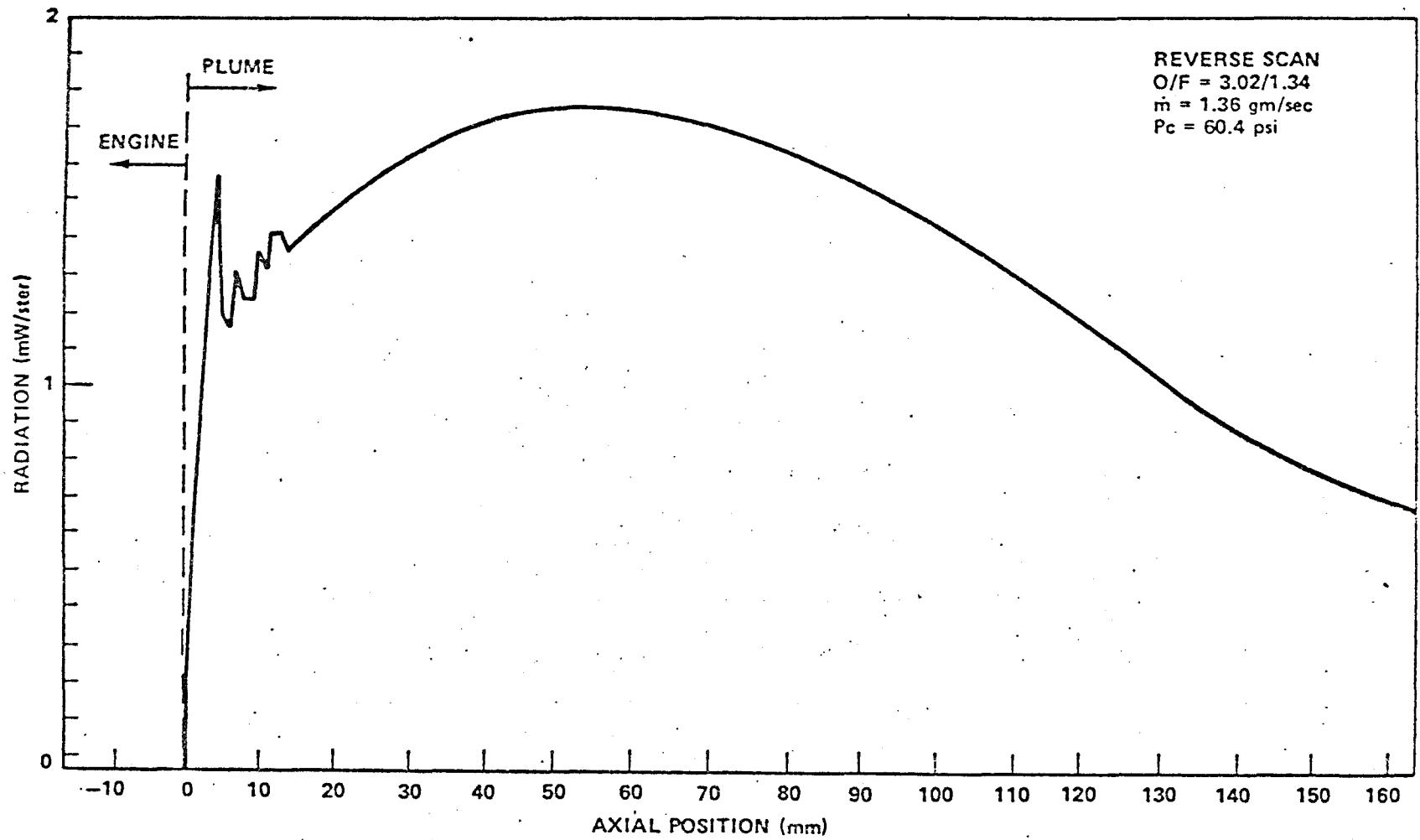


Fig. 3.22 Axial variation of radiation at O/F ratio of 2.27.

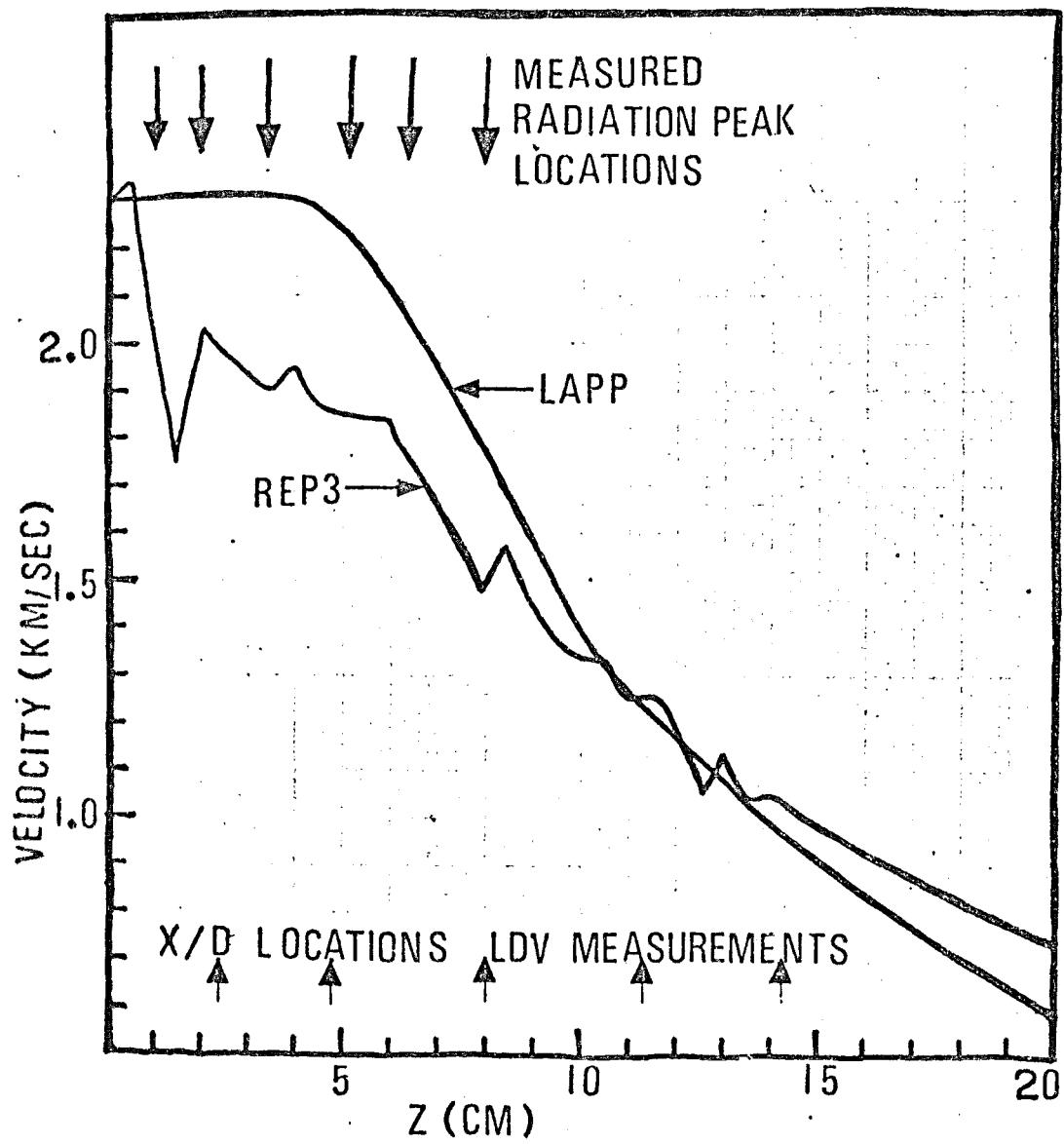


Fig. 3.23 Calculated axial velocity distribution.

several other oscillations in velocity, as shown by the jagged nature of the curve. A particularly large decrease in velocity occurs near an X/D value of 8. Also noted in Fig. 3.23 are the locations of the measured radiation peaks shown previously in Fig. 3.20. The X/D locations for the LDV measurements also are indicated on Fig. 3.23.

The predicted variations of temperature at the jet centerline with axial distance are given in Fig. 3.24. The LAPP code predicts a small temperature rise due to afterburning in the exhaust plume. The TKE code predicts a sharp rise in temperature at an X/D value of one, which then is followed by a series of oscillations in temperature. The predicted oscillations do not agree with the measured locations of radiation peaks (denoted by arrows on the Figs.). It is quite obvious that neither of the computer codes does an adequate job of predicting the spatial variation of velocity and temperature produced by shock waves and Mach disks in the near field of the exhaust plume. This is not surprising since LAPP has no provision for treating shocks.

Figures 3.25-3.38 present the variation of the axial component of the velocity with radial distance from the jet centerline. As shown in Fig. 3.25 for an X/D value of 2.4, the predicted results (curves) do not agree with the experimental measurements (symbols). Because of the limited number of particles in the hot core of the jet, the NASA LDV system was not able to make measurements closer to the jet centerline than about one exit nozzle radius. The variation of the radial velocity component also is shown. The TKE code predicts an increasing value of v with radius, whereas the measured value decreases with radial distance. The external flow is an axial flow at 61 m/sec.

The lack of agreement between theory and experiment also is indicated in Fig. 3.26 for the velocity profile at an X/D value of 8.0. The LAPP and TKE

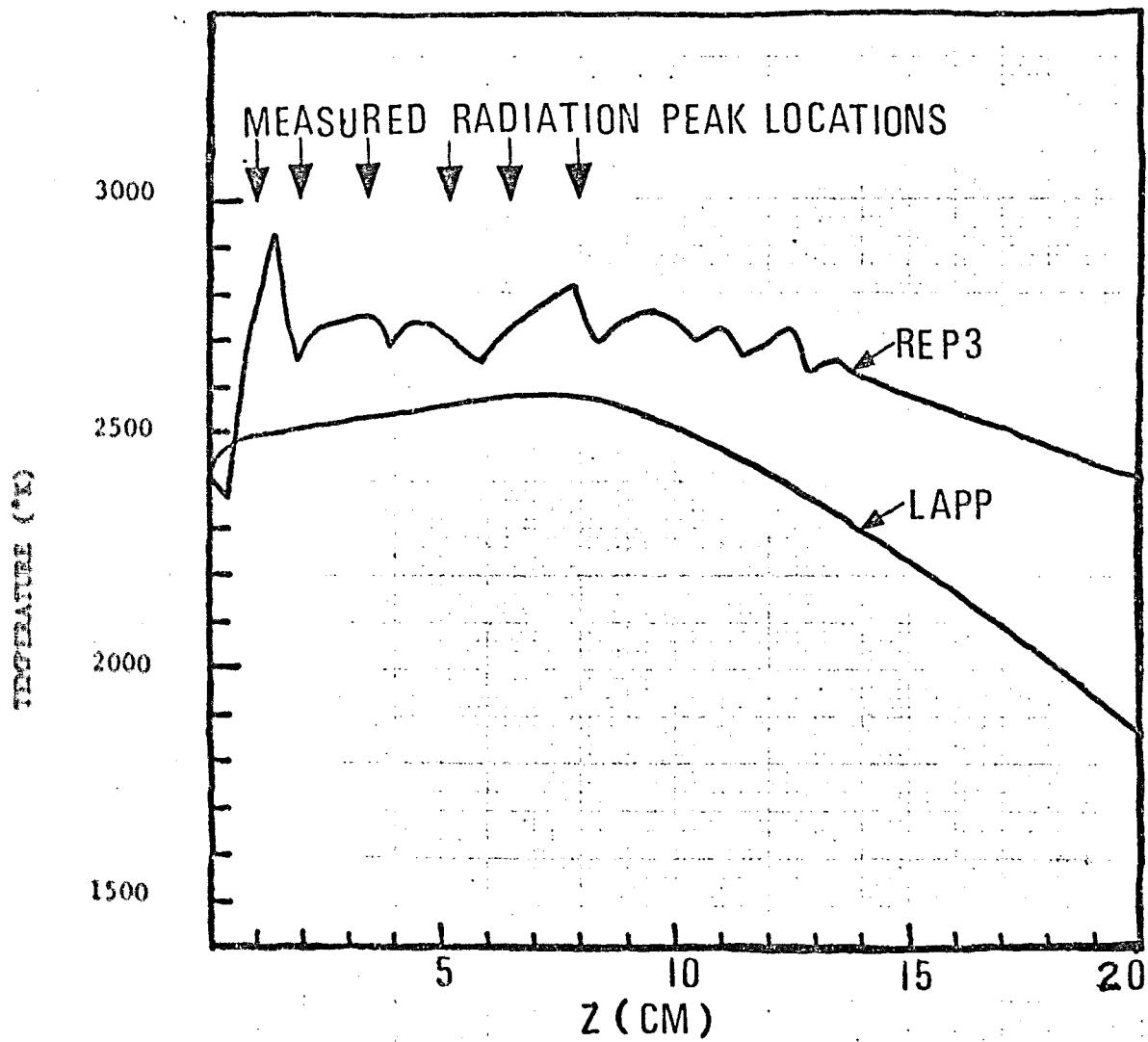


Fig. 3.24 Calculated temperature distribution.

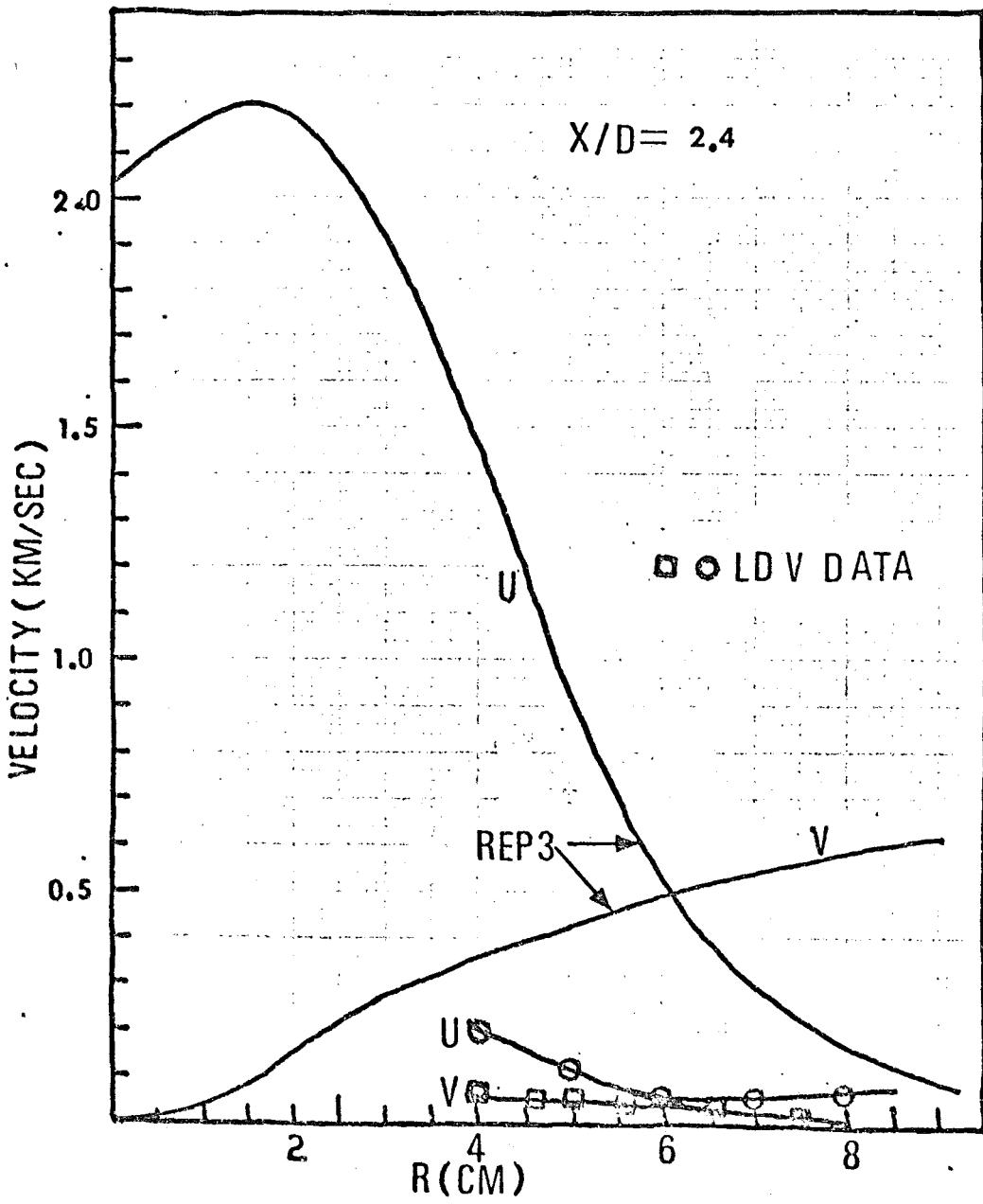


Fig. 3.25 Velocity profiles at $X/D = 2.4$

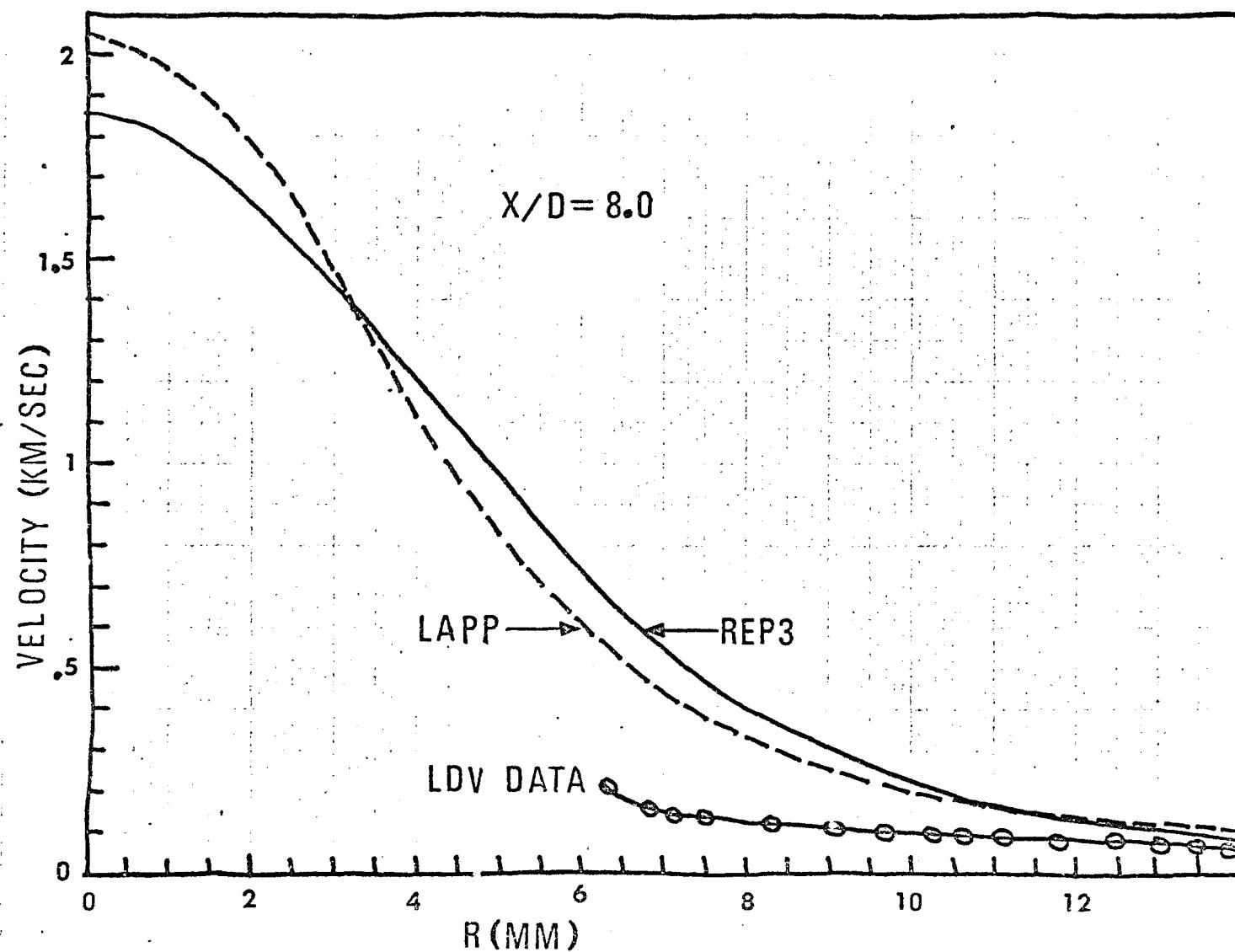


Fig. 3.26 Velocity profiles at $X/D = 8.0$

calculated results are similar but are in excess of the experimental values.

Similar results are presented in Figs. 3.27 and 3.28 for X/D values of 11.3 and 14.2.

The LDV data indicate that the exhaust plume is much narrower than predicted by the LAPP and TKE gas dynamics models. Since both the LAPP and TKE models assumed that turbulent mixing began at the nozzle exit plane and ignored the large recirculation (base flow) region, it is expected that turbulent mixing occurs too fast compared to the physical situation where turbulent mixing begins at the end of the recirculation zone, which is several nozzle diameters downstream from the exit. It is known that the base flow region can have a pronounced effect on exhaust plume spread and initial turbulent mixing layer growth. Because of time and budget constraints, we were unable to use a base flow model to assess its effect on the calculated temperature and velocity profiles.

The infrared radiation band model program was used to calculate the variation of infrared radiation along the exhaust plume for the temperature and species concentration profiles predicted by the LAPP and TKE gas dynamics computer codes. The computed results are shown in Fig. 3.29. The lack of agreement between theory and experiments was not unexpected, since the calculated temperature distributions did not yield the spatial structure due to shock waves that are evident in the measured infrared data and are visible in photographs of the exhaust plume.

It is not clear what causes the difference between the REP3 and LAPP calculations. The LAPP code plus the radiation code yield results that are within a factor of 1.5 of the measured results while the factor is approximately 2.5 for the REP3 plus radiation code results.

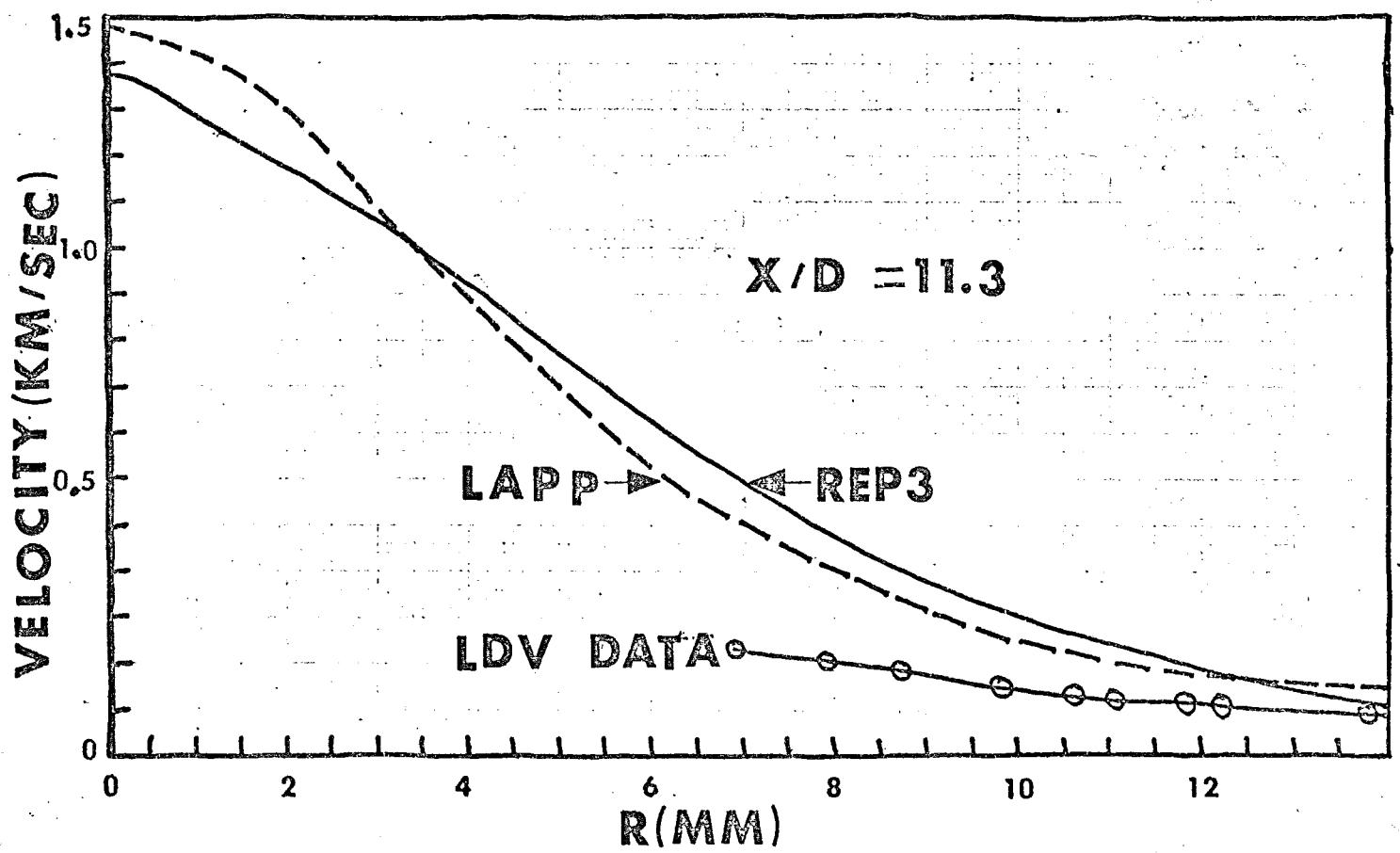


Fig. 3.27 Velocity profiles at $X/D = 11.3$

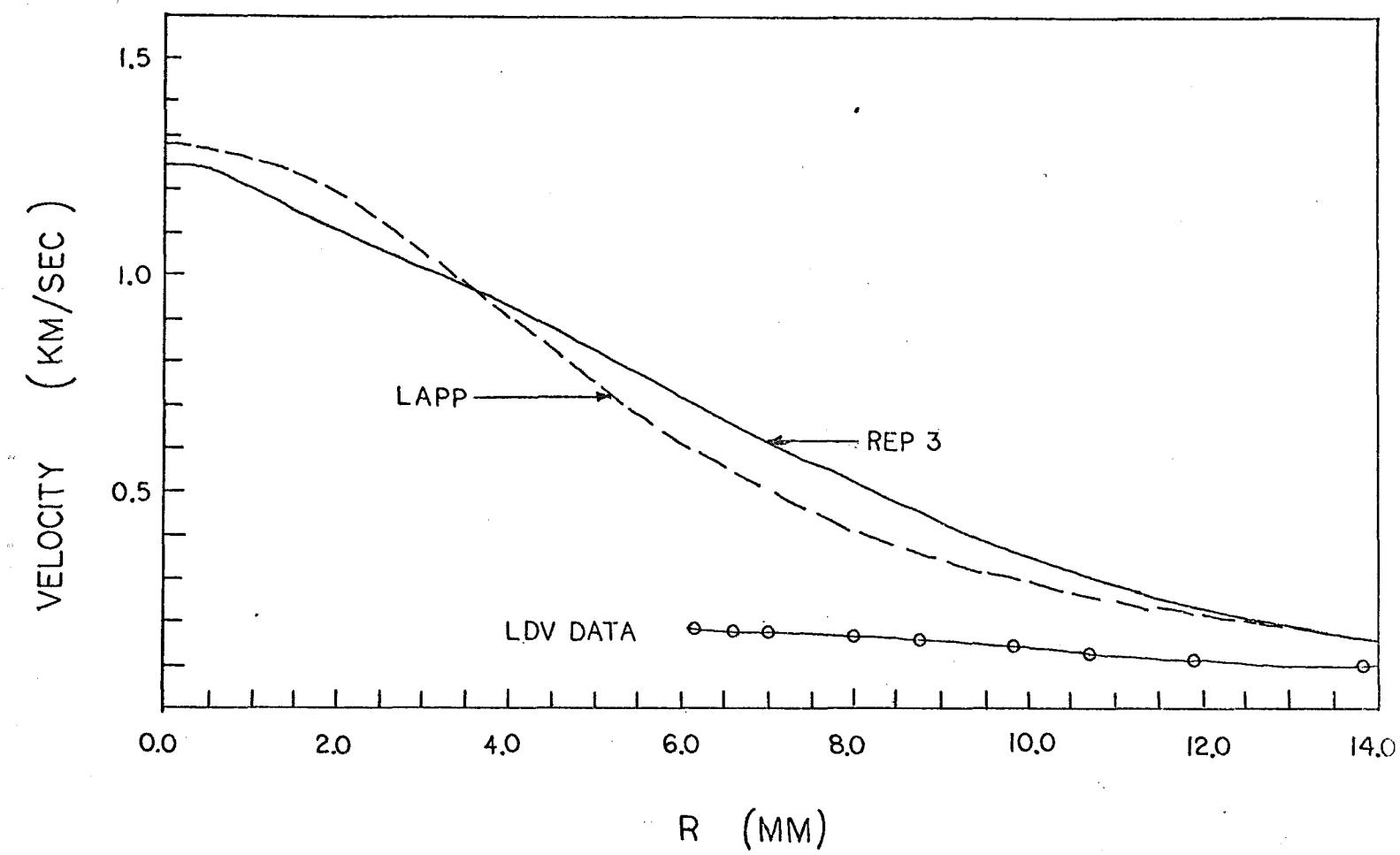


Fig. 3.28 VELOCITY PROFILES AT $X/D = 14.2$

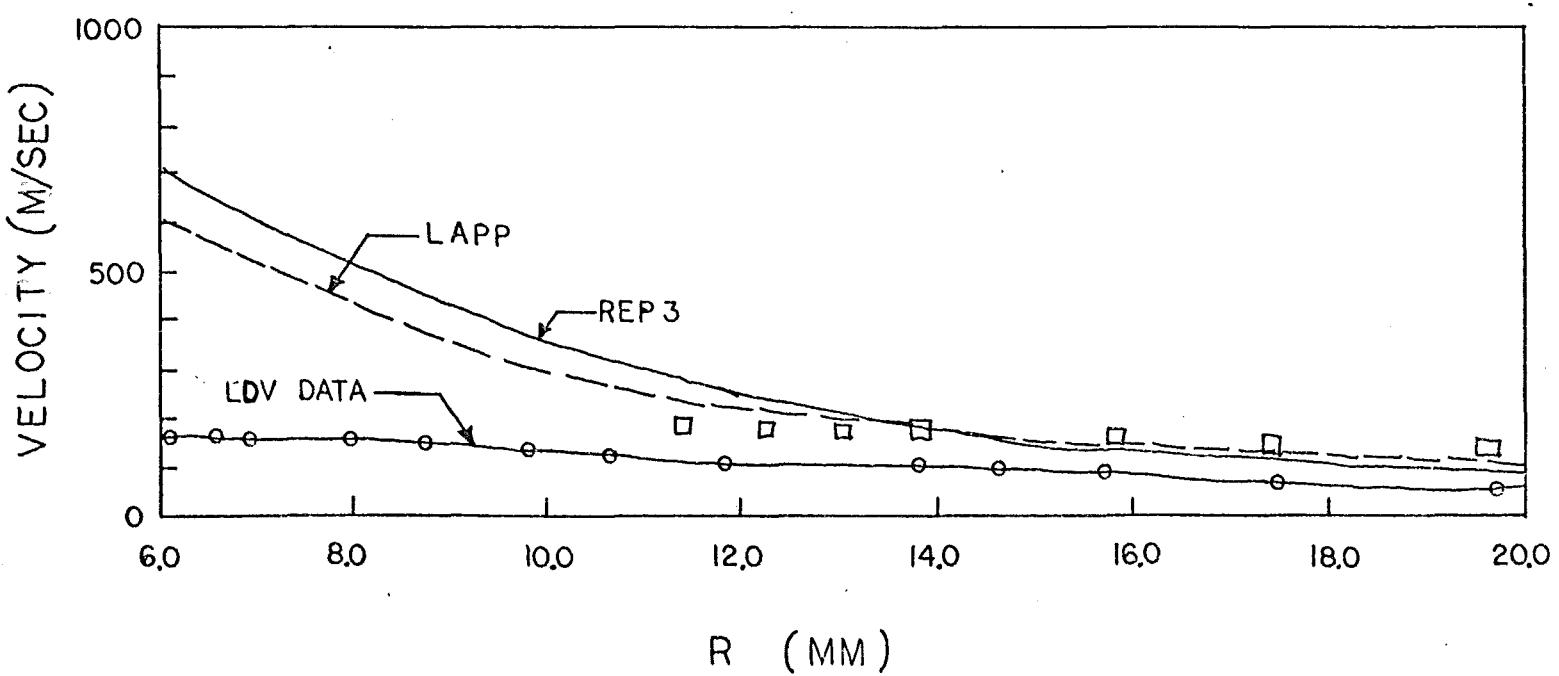


Fig. 3.28 (continued)

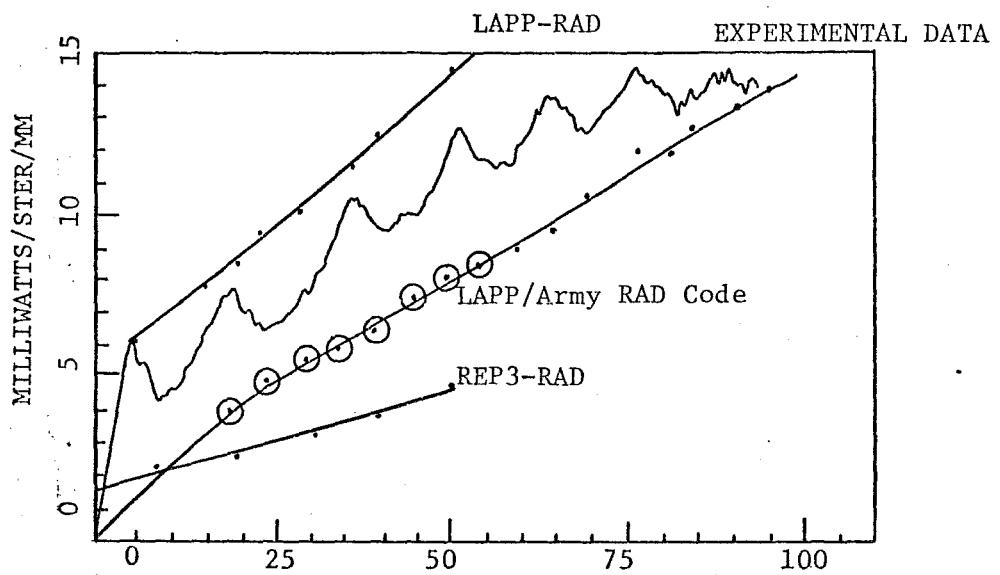


Fig. 3.29 Comparison of calculated and measured IR radiation.

3.4 Gas Dynamics Finite Difference Computer Model

A Gas Dynamics Finite Difference Computer Model was developed as part of the contract effort. Numerous difficulties and problems were encountered during the development process, but the program was operational at the conclusion of the effort. Unfortunately, the present version of the code is not entirely satisfactory due to the fact that the computational times are large. Approximately one hour on the IBM 360 computer system was required to march the solution 360 cycles (which corresponds to a flow time of 0.2 msec.). While this time would be approximately 6-10 minutes on the CDC 660 computer, it is still felt to be quite excessive since the test flow being analyzed was a relatively simple expanding flow. In its present form, the finite difference code successfully calculates shock waves on wedges and blunt bodies and is moderately successful in calculating the shock-Mach disc structure in a cold gas nozzle expansion.

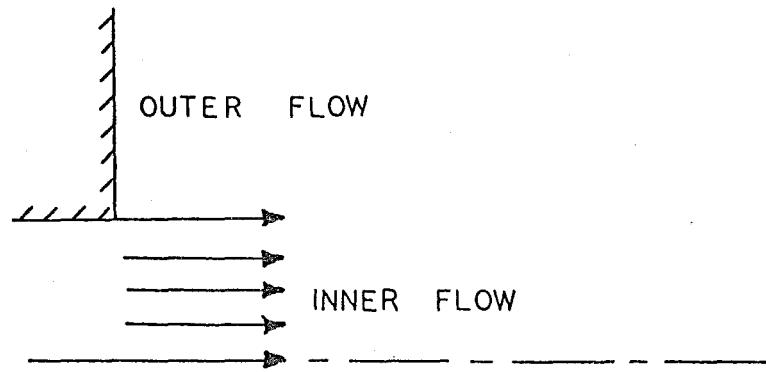
Several typical nozzle expansion test cases were computed to assess whether the numerical results were reasonable. These results are described in the next section.

3.4.1 Overexpanded Nozzle Exhaust

In order to assess the FD model's ability to handle strong shock waves which result in the exhaust plume from an overexpanded nozzle, the flow having the initial characteristics given in Fig. 3.30.

The computed Mach number variation along the jet centerline is given in Fig. 3.31. The Mach number decreased to a subsonic value in a very short distance downstream from the exit plane and remained subsonic for 80-cm in the flow direction.

An analytical estimate of the flow conditions was made by assuming that an oblique shock wave would be formed at the nozzle exit and would raise the



INNER FLOW

$M_{exit} = 1.133$

$T_{exit} = 300^{\circ}\text{K}$

$P_{exit} = 0.8 \text{ atmospheres}$

OUTER FLOW

$M = 0$

$T = 300^{\circ}\text{K}$

$P = 1.0 \text{ atmospheres}$

Fig. 3.3.0 INPUT FLOW PARAMETERS

MACH NUMBER

1.1

1.0

.9

.8

.7

.6

$Z = .05$

Distance along centerline (cm)

CALCULATED OBLIQUE SHOCK AT EXIT (SEE ATTACHED)

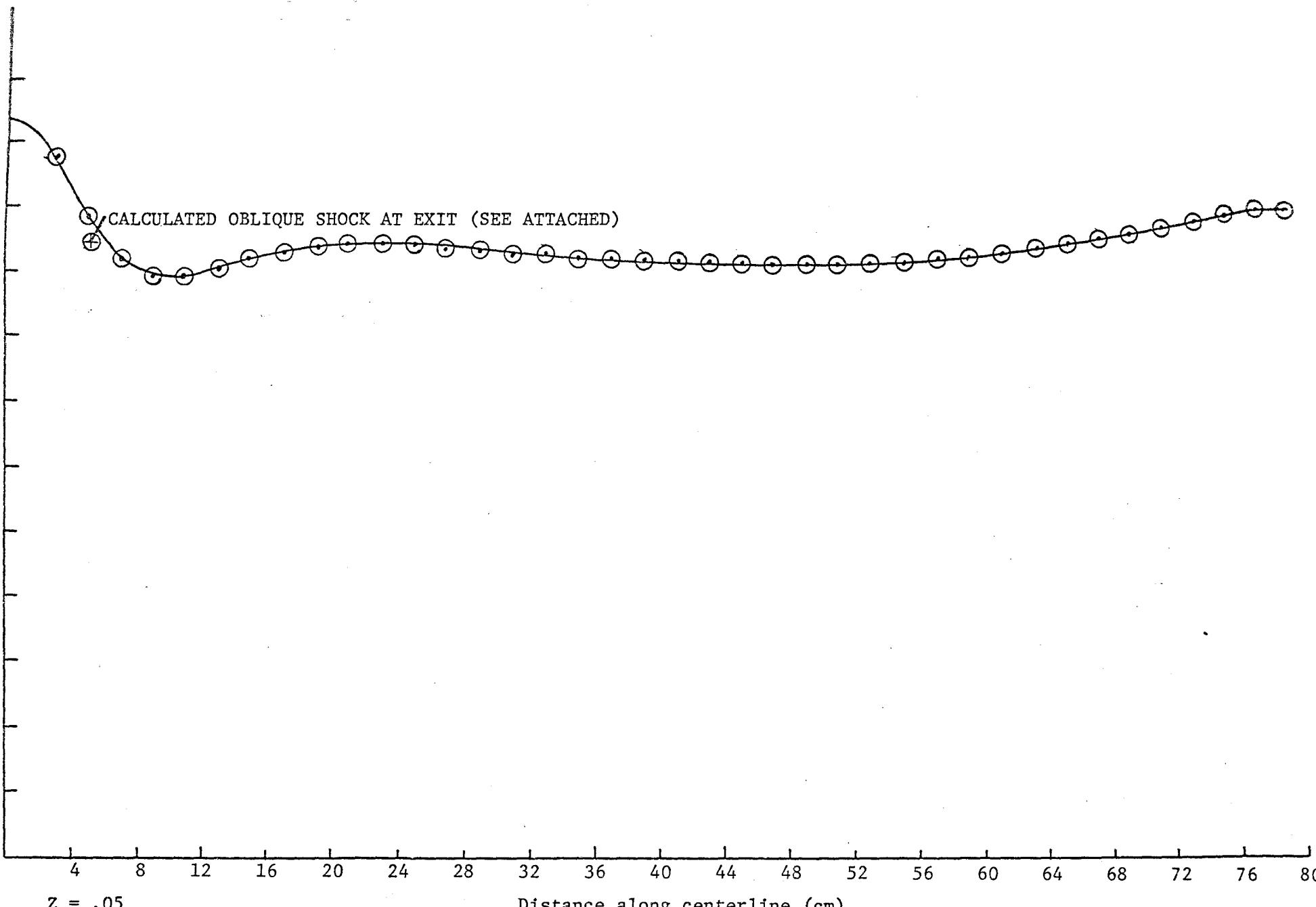


Fig. 3.3.1 CENTERLINE MACH NUMBER VARIATION

pressure behind the shock wave to the ambient pressure of 1.0 atm. The oblique shock angle was then determined to be approximately 76.6° as shown in Fig. 3.32. The shock wave would intersect the centerline axis at a distance of 0.2323-cm downstream of the nozzle exit plane. Gas dynamics calculations yielded a value of 0.94 for the Mach number in region (2) behind the shock wave. This calculated value is indicated on Fig. 3.31 and is in good agreement with the finite difference code result.

As shown in Fig. 3.32, the shock wave would reflect from the centerline and would process the flow back parallel to the centerline axis. In doing so, the pressure in region (3) would increase while the Mach number M_3 would be less than M_2 . This agrees qualitatively with the FD code results.

At the intersection of the reflected shock with the free boundary, an expansion wave would be generated to reduce the pressure at the edge of the plume back to one atm. The expansion fan would accelerate and expand the flow in regions (4-6) until a new shock pattern would be formed between regions (7 and 8). Then the entire pattern would repeat.

While the FD code calculations appear to be quantitatively correct, there were no available numerical data to compare with the FD code results in the subsonic flow region. In order to determine whether the FD code predictions are reasonable in an expanding flow, another test case was designed in which a supersonic flow would expand from a nozzle. The results for this flow are described in the next section.

3.4.2 Underexpanded Nozzle Exhaust

As stated above, a test case involving the supersonic expansion of air from a nozzle was employed so that the FD code predictions could be compared with the results obtained using the Method of Characteristics.

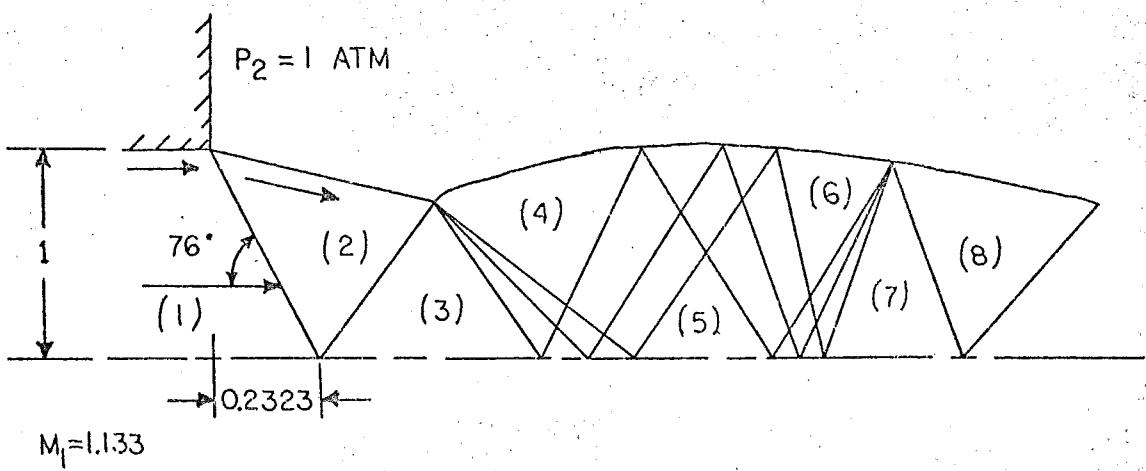


Fig. 3.32 ASSUMED SHOCK PATTERN FOR HAND CALCULATION.

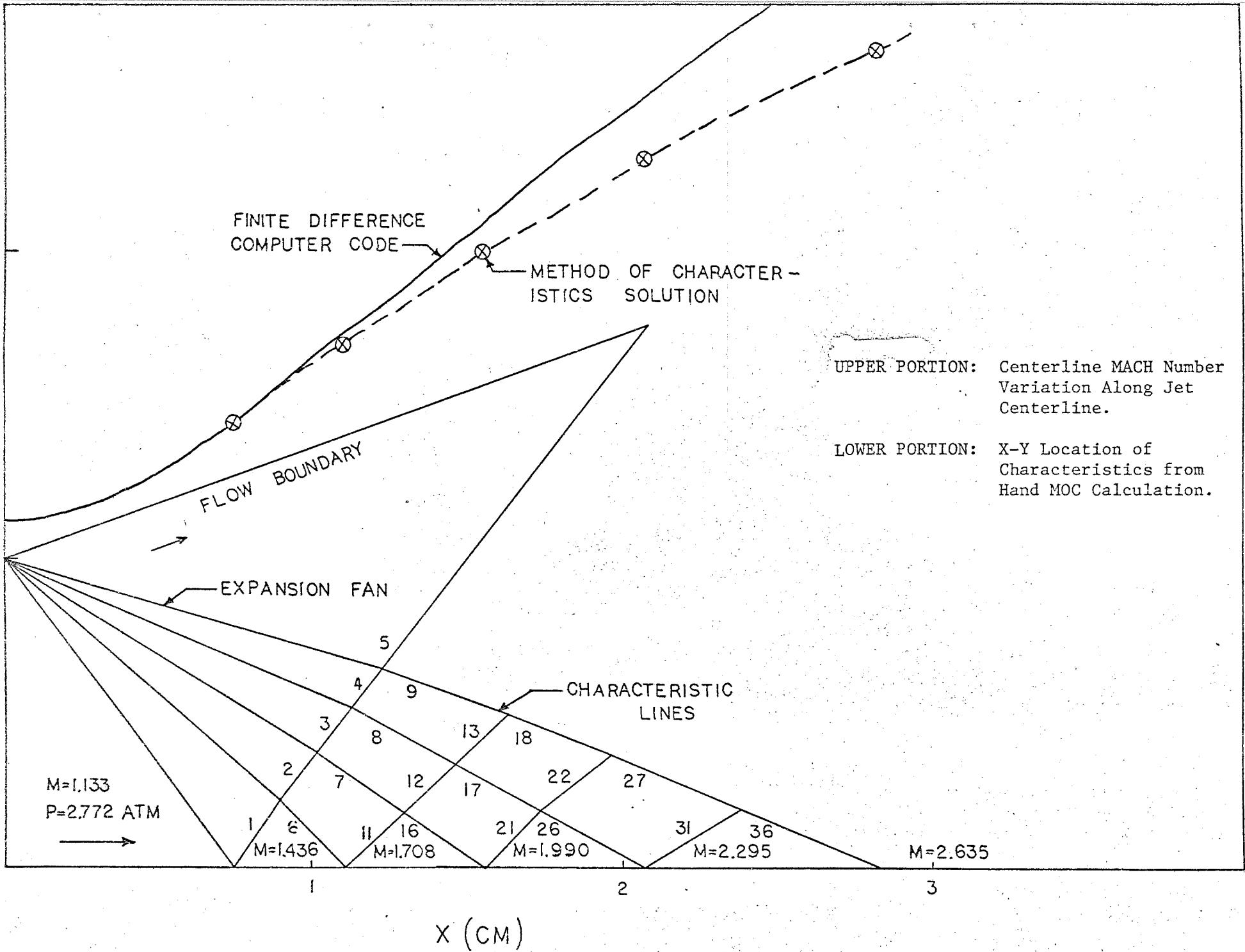


Fig. 3.33 COMPARISON OF METHOD OF CHARACTERISTICS AND FINITE DIFFERENCE MODEL RESULTS

The method of Characteristic technique used was Temple's Method as taught by the Principal Investigator in a first-year graduate course in gas dynamics. The characteristics net produced in the X-Y physical plane is shown in the lower portion of Fig. 3.33. A comparison of the MOC results (using 5 characteristics to start the calculation) and the Finite Difference Code results is presented in the upper portion of Fig. 3.33. The results are in good agreement. The lack of agreement at large distances from the nozzle exit plane is to be expected since the hand-calculated MOC solution used only five characteristics in the initial fan and is therefore expected to be more approximate than the Finite Difference Computer Code which uses a more complete set of differential equations.

Even though we were encouraged that the computer code appears to yield reasonable numerical results for an expanding plume, we are discouraged by the fact that the computer program required approximately one hour on the IBM 360 computer system to march the solution 0.3 msec. (or 360 cycles). While this time would be approximately 6-10 minutes on the CEC 6600, it still is felt to be quite excessive in view of the simple flow calculated.

3.4.3 Other Test Cases

In earlier work [21], the finite difference code was used to calculate the viscous flow between two parallel walls with one of the walls moving relative to the other. This is the so-called classical Couette Flow problem whose theoretical solution is well known [32]. The computed numerical results agreed well with the theoretical model results. In addition, the computer solution converged in 10 to 20 iteration cycles and the total computational time on the IBM 370 computer was approximately 30 seconds.

The code also successfully predicted the flow over a two-dimensional wedge in which a strong bow shock wave is formed.

In both of these cases, we have solid boundaries which restrict the flow and appear to assist in the convergence of the solution to a stable value.

3.4.4 Chemical Reactions and Turbulent Mixing Computations

Due to problems with the basic finite difference computer code, the chemical reaction and turbulent mixing portions of the code were not fully operational at the end of the program.

3.5 Results Obtained Using the REP 3 TKE Computer Model

The REP3 TKE Computer Model described earlier was revised for use on the CDC computer at MIRADCOM. The test case supplied by Dr. David Jensen of the Rocket Propulsion Establishment of Great Britain was successfully run on the MIRADCOM computer.

After the REP3 program was operational, it was utilized to calculate the flow field from a small kerosene/gaseous oxygen engine similar to the one used as the test case by Jensen. The results of these calculations are presented in Figs. 3.23-3.28. As discussed in Section 3.3, the REP 3 predictions do not agree with the experimental data. Attempts to locate any problems with the calculations or with the experimental data were unsuccessful.

One computer run was devoted to rerunning the case described in the previous paragraph. The program was rerun using a much smaller grid spacing. The results are shown in Fig. 3.34. The top curve shows the oscillations in temperature while the lower curve indicates the results for the velocity. Also shown on the graph are the locations of the radiation minima (two experimental runs are shown: Runs 65 and 85). As can be seen, the first three locations of the calculated temperature minima agree with the experimental locations. Even though the fourth through seventh minima do not agree, the results are reasonable.

ARROWS LOCATE RADIATION PEAK LOCATIONS

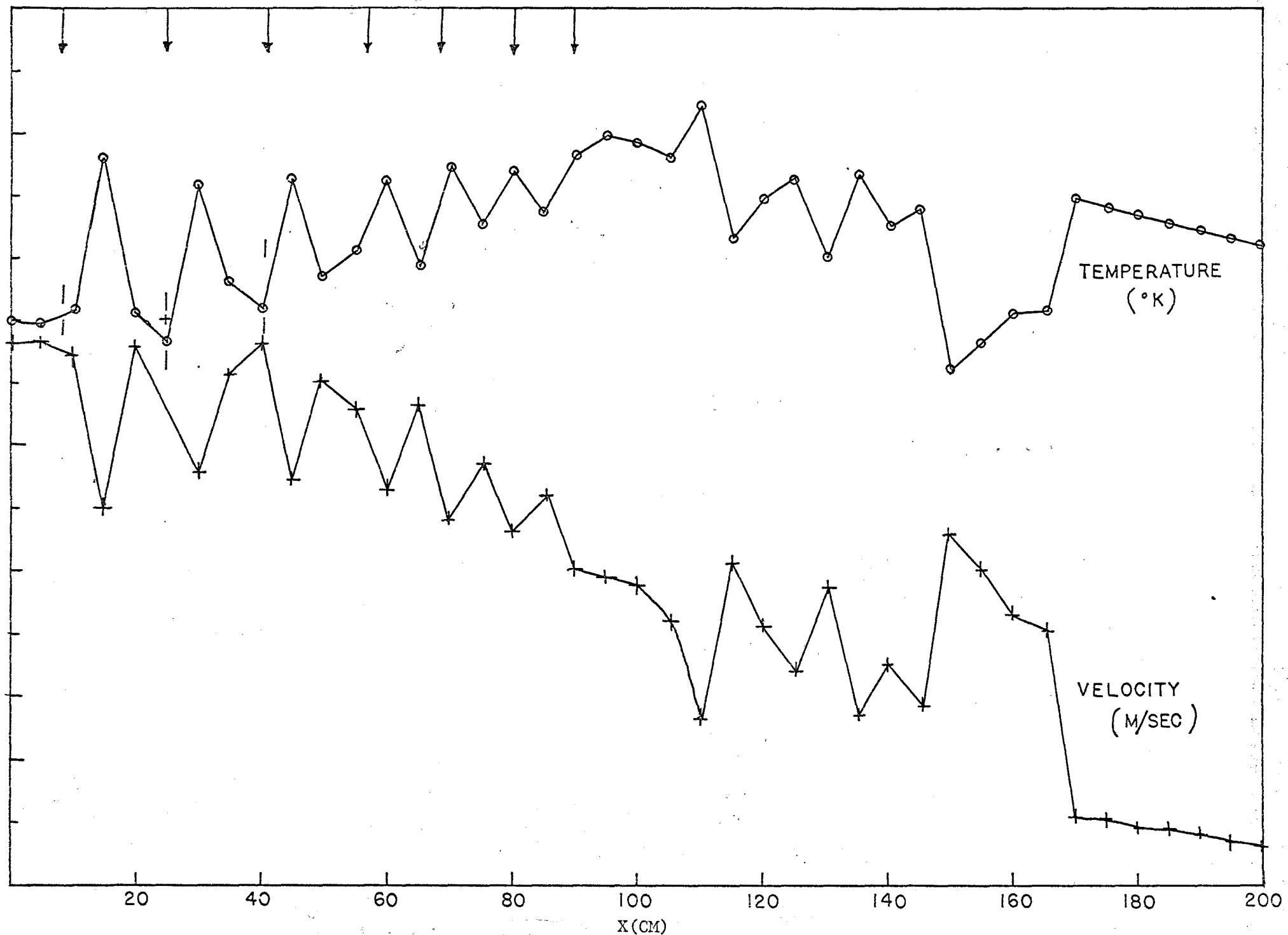


Fig. 3.34 CALCULATED RESULTS USING THE REP3 COMPUTER CODE

However, the behavior of the calculated profiles for X greater than 90-mm are disturbing in that there are a series of temperature and velocity oscillations which do not agree with any experimental or physical evidence. Furthermore, the predicted sharp drop (rise) in temperature (velocity) at 145-mm and the sharp rise (drop) in temperature (velocity) at 165-mm do not correspond to observed physical phenomena.

Clearly, further efforts are needed to ascertain whether or not the REP3 Program is correct. Unfortunately, there was no available time or funds to consider this further.

It was also discouraging that over 30 minutes of computer time on the CDC 6600 was required to obtain the small step size required to resolve the near field shock waves.

4. SUMMARY AND CONCLUSIONS

A research project was conducted which resulted in the development of an extended version of the MIRADCOM (Jackson) Infrared Radiation Computer Model.

The extended computer code was validated by comparing the predicted spectral radiance of H_2O and CO_2 in the $2.7-\mu m$ and $4.3-\mu m$ wavelength bands with experimental data and with predictions of the NASA (Reardon) Radiation Computer Code. The predicted results agree fairly well with the NASA code even though there were some slight differences. With the exception of a few results for low pressure CO_2 , the predicted values were within $\pm 20\%$ of the measured values.

The infrared computer code was used to predict the infrared radiation emitted in the wavelength interval 4 to 5 microns from the exhaust of a full scale turbojet aircraft. The predicted results appeared to be reasonable when compared with the available test data.

The infrared radiation band model was used to calculate the variation of infrared radiation along the exhaust plume of a small kerosene/gaseous oxygen rocket engine. The temperature and species concentration profiles were predicted using the LAPP and REP3 computer codes. The agreement between the theoretical calculations and the experimental measurements was poor. The lack of agreement was due to the fact that the calculated temperature distributions did not exhibit the spatial structure due to shock waves that are evident in the measured infrared data and are visible in photographs of the exhaust plume. There appears to be short comings in the ability of both the LAPP and REP3 codes to predict the correct spatial distributions of temperature and specie concentrations.

The Gas Dynamics Finite Difference Computer Model developed as part of this project was only partially completed and was able to produce reasonable numerical results at the expense of long computer runs and high computer memory requirements. The code successfully calculated the flow field for flow geometries having solid boundaries, but had problems with the free boundaries encountered in nozzle exhaust plumes. The results obtained for the near field of the nozzle appeared to be correct, but the downstream results were not validated.

The preliminary calculations obtained using the REP3 and TKE computer codes did not agree with the experimental data.

Due to the partial success of the finite difference and REP3 computer codes, further development and assessment are required before these codes are fully operational in a validated model.

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APPENDIX I

LISTING OF IR COMPUTER PROGRAM

IIXIXII II I I II IIXIXII

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//HARWE JOB (36931,51118,REMOTE3,R-192,T-35,U-10),NGUYEN  
//STEP1 EXEC FORTGCLG,FORTREG=192K,GOREG=192K  
//FORT.SYSPRINT DD DUMMY  
//FORT.SYSIN DD *  
COMMON /PLOT1/WAVEI,WAVEF,CENTR,TOTRAD  
COMMON /RASP/RASP  
COMMON /BANDP/ XKK(254,7),DDINV(254,7),UK(254),AKU(254)  
DIMENSION I6(3),I8(10),I7(6),I2(10),I3(6),I4(4),I5(10),I1(5)  
COMMON /RPN1/ NWAV,WAVNO(254),WAVLT(254)  
COMMON /CHECK/IREAD  
COMMON /INDATA/ X(40),R(40,30),T(40,30),PCO2(40,30),  
$ PH20(40,30),PN2(40,30),PO2(40,30),PCO(40,30),PH2(40,30)  
COMMON /CTRL/ TITDE(12),KOUNT,NXS,NORUN  
COMMON /OTH/ PP,RJ,NRP(40)  
COMMON /SWITCH/ KSWIT ,IXXJG
```

C

```
DIMENSION RADSTL(254),RADXT(254),TRILAST(254) 000130  
DIMENSION CR(100),CX(100),CL(100) 000140  
DIMENSION RADXI(254),RRN(100)  
DIMENSION STRAD(40)  
DIMENSION RAD(254),RADX(254),TRNS(254),P(254),O(254),RADSTR(254) 000230  
DIMENSION WAVL5(254) 000240  
DIMENSION RCALC(100) 000260  
DIMENSION SRAD(254),SRADSR(254)  
INTEGER IH20(2),ICD2(2),IWRT(2)  
REAL HRADXT(40),CRADXT(40)  
REAL HCTOTX(40),HCRDXT(40)  
REAL HRADX(254),CRADX(254),SRADX(254)  
EQUIVALENCE (HRADX(1),RADX(1)),(CRADX(1),RAD(1)),  
$ (SRADX(1),TRNS(1))  
EQUIVALENCE (HCTOTX(1),CR(1)),(HCRDXT(1),CX(1))
```

C

```
DATA I6N/10/  
DATA I8N/37/  
DATA I7N/23/  
DATA I2N/33/  
DATA I3N/25/  
DATA I4N/19/  
DATA I5N/49/  
DATA I1N/20/  
DATA ASTER,RCALC/1H*,100*1H / 000300  
DATA IDISC/8/  
DATA ISTOP/0/  
DATA IDSCT/0/  
DATA JPASS/0/  
DATA IFLG/0/  
ISWIT = 2  
KSWIT = 2  
JSWIT=2  
LSWIT = 1  
IREAD=0  
KPZQ=0  
VZU=1  
KNPQ=0  
IXXJG=0
```

C

```
C SENSE SWITCH 1 SET--YES,PRINT FLOW FIELD DATA  
C SENSE SWITCH 2 SET--YES,CO2 RADIATION--NO,H2O RADIATION  
C SENSE SWITCH 3 SET--YES,AEROCHEM INPUT TAPE--NO, MOTOR DATA INPUT TAPE  
C SENSE SWITCH 4 SET--YES,H2O AND CO2 RADIATION COMBINED
```

TOTRAD=0.0

000590

C
C SET PRINT LABEL
C
C DETERMINE MODE OF INPUT TAPE -- MOTOR DATA OR AEROCHEM
C
172 CONTINUE
END FILE 9
REWIND 9
END FILE 8
REWIND 8
IF(KPZQ.EQ.1) GO TO 104
GO TO (101,102),JSWIT
102 CONTINUE
CALL READER
KPZQ=1
GO TO 104
101 CONTINUE
104 CONTINUE
C
NXS=KOUNT
DO 111 I=1,NXS
IF(NRP(I).GT.25) GO TO 111
NRP1=NRP(I)+1
ISET=NRP(I)
DO 112 KK=NRP1,25
R(I,KK)=R(I,ISET)
T(I,KK)=T(I,ISET)
PCO2(I,KK)=PCO2(I,ISET)
PH2O(I,KK)=PH2O(I,ISET)
PN2(I,KK)=PN2(I,ISET)
PO2(I,KK)=PO2(I,ISET)
112 CONTINUE
111 CONTINUE
C
C READ WAVELENGTHS AND CORRESPONDING BAND PARAMETERS
C READ IN H2O BAND DATA FIRST FOLLOWED BY CO2 BAND DATA
C
47 CONTINUE
IF(LSWIT .EQ.1.AND.ISTOP.EQ.3.AND.IFLG.LT.2) GO TO 21
IF(ISTOP.EQ.3) GO TO 23
21 CONTINUE
WRITE (6,9911) 000330
READ (5,9912) NWAV1, NWAV2 000340
WRITE (6,10035) NWAV1, NWAV2
10035 FORMAT(10H NWAV1 ,2(1X,I5))
9912 FORMAT(2I10)
NWAV1=1 + ((NWAV2-NWAV1)/25) 000360
DO 1 I=1,NWAV 000380
IF(KNTO.EQ.0) GO TO 9957
GO TO 9967
9957 READ (5,9922) WAVN7(I),(XKK(I,J),J=1,7)
9922 FORMAT (F9.1,7E9.2) 000400
GO TO 9977
9967 READ (5,9922) WAVN7(I),(XKK(I,J),J=1,7)
9977 CONTINUE
WRITE (6,10037) WAVN7(I),(XKK(I,J),J=1,7)
10037 FORMAT(10H WAV ,F9.1,7E9.2)
IF(KNTO.EQ.0) GO TO 9987
GO TO 9997
9987 READ (5,9001) WAVN7(I),(DDINV(I,J),J=1,7)
9001 FORMAT(F6.1,E9.2,6E10.2)
GO TO 9947
9997 READ (5,9010) WAVN7(I),(DDINV(I,J),J=1,7)
9010 FORMAT (F6.1,E9.2,6E10.2)

```

9947 CONTINUE
  WRITE (6,10037) WAVNO(I),(DDINV(I,J),J=1,7)          000420
  WAVLT(I)=10000./WAVNO(I)
9910 FORMAT (1X,8E15.7)                                000450
  1 WAVLS(I)=WAVLT(I)**5
  WAVEI=WAVNO(1)                                       000460
  WAVEF=WAVNO(NWAV)                                    000470
  WRITE (6,9911)                                         000480
9911 FORMAT (1H1)
  DO 9913 I=1,NWAV                                     000490
  DO 9913 J=1,7
    IF (DDINV(I,J).LE. 0.0) DDINV(I,J)=1.0E-20        000500
    IF (XKK(I,J).LE.0.) XKK(I,J)=1.0E-20
9913 CONTINUE                                           000510
9000 FORMAT(F9.4,7E9.3)                                000520
000530
C
C   PRINT BAND PARAMETERS IF OPTIONED SELECTED
C
IF(LSWIT .EQ.1.AND.ISTOP.EQ.0.AND.IFLG.LT.1)CALL SAVE1
IF(ISWIT .EQ.1) CALL BPPRT
C
23 CONTINUE
IF(LSWIT .EQ.1.AND.ISTOP.EQ.3.AND.IFLG.GE.2) CALL SWAP
IFLG=IFLG+1
C
C   DETERMINE INPUT DATA PRINT OPTIONS
C
C   SUBROUTINE PRMD WRITES OUT INPUT FLOW FIELD DATA
IF(ISWIT .EQ.1) CALL PRMD
C
C   ALPHA=ASPECT ANGLE IN DEGREES(0 DEG=NOSE-ON)      001520
C   ISTOP=1,WILL TERMINATE PROGRAM AFTER SINGLE ALPHA CALCULATION 001530
C   ISTOP=2,RETURNS TO THIS POINT FOR ANOTHER ALPHA       001540
C   ISTOP=3,RETURNS TO STMT 172 TO READ NEW SET FLOW FIELD DATA 001550
C   SENSE SWITCH 4 WILL OVER RIDE ANY ISTOP OPTION       001560
C
103 READ(5,1010)ALPHA,ISTOP                           001570
1010 FORMAT(F10.5,T5)
  ALP=ALPHA/57.295780                               001580
  CC=COS(ALP)                                       001590
  SC=SIN(ALP)                                       001600
  PI=3.1415926                                     001610
  TRADX=0.0                                         001620
  TOTRAD=0.0                                         001630
  CENTER=0.0                                         001640
  C
  DO 185 IL=1,NWAV                                  001650
185 RADXT(IL)=0.0                                     001660
C
  WRITE (6,1401)
  DO 3000 I=1,NXS                                 001670
  NZPEX=NRP(I)
  WRITE (6,1401)
1401 FORMAT (1H1)
  WRITE (6,1404) (TITLE(JRR),JRR=1,12)            001680
  001690
1404 FORMAT(10X,12A4,A2,/)

  WRITE (6,1400) T,X(I),ALPHA
1400 FORMAT (1X,12)STATIN N0.=,14,5X,17HAXIAL DIST. (CM)=,F10.3,5X,
$18HASPECT ANGLE (DEG)=,F5.1,/
  TRADX=0.0                                         001700
  IXD=1                                             001710
  RD=R(TXD,25)
  IF (RD .LE. 0.0) RD=1.0E-20                     001720
  X=I=X(IXD)+.01
  IF(ALPHA.EQ.180.0) X=0(XNS)
  ZPEX=IZPEX                                         001730
  001740
  001750
  001760
  001770
  001780
  001790
  001800
  001810
  001820
  001830
  001840
  001850

```

```

C ***CHANGE DELZ=R0/ZPEZ TO DELZ=1. FOR THE SLAP***  

DELZ=1.0  

DELZ=R0/ZPEX  

C ***** SET NZU=NZPEX FOR TOTAL PLUME *****  

ZU=RASP/DELZ  

NZU=NZPEX  

C  

DO 190 II=1,NWAV 001870  

SRADSR(II)=0.0  

190 RADX(II)=0.0 001880  

STOT=0.  

C  

C THIS CARD HAS BEEN CHANGED AS  

C DO 2000 M=1,NZPEX  

DO 2000 M=1,NZU  

ZM=M 001910  

ZD=DELZ*(ZM-1.) 001920  

IF(ALPHA.EQ.0.0) RD=2.*R(1,25)+(ZM-1.)*(R(NXS,25)-2.*R(1,25))/ZPEX 001930  

IF(ALPHA.EQ.180.0) RD=(ZM-1.)*(R(NXS,25))/ ZPEX 001940  

C 001950  

C***** 001960  

C CALCULATE COORDINATES X,R ALONG LINE-OF-SIGHT AS FUNCTION OF ITS 001970  

C LENGTH 001980  

C 001990  

DELL=10. 002000  

IF(ALPHA.EQ.0.0.OR.ALPHA.EQ.180.0) DELL=X(NXS)/98. 002010  

200 XL=0. 002020  

C  

DO 203 J=1,100 002030  

N=J 002040  

APR=XL*SC-SQRT(R0*RO-ZD*ZD) 002050  

CR(N)=SQRT(ZD*ZD+APR*APR) 002060  

CX(N)=X0+XL*CC 002070  

CL(N)=XL 002080  

IF(CX(N).LT.0.) GO TO 204 002090  

IF(CX(N).GT.X(NXS)) GO TO 204 002100  

IF(ALPHA.EQ.0.0.OR.ALPHA.EQ.180.0) GO TO 203 002110  

DO 201 K=1,NXS 002120  

KI=K 002130  

IF(CX(N).LE.X(K)) GO TO 202 002140  

201 CONTINUE 002150  

202 IF(CR(N).GT.R(KI,25) .AND. N .GT. 10) GO TO 204 002170  

203 XL=XL+DELL 002180  

C 002190  

C DELL TOO SMALL, INCREASE AS FOLLOWS 002200  

C 002210  

DELL=DELL*1.5  

GO TO 200  

204 CONTINUE  

IF(N.GE.50) GO TO 299 002260  

C 002270  

C DELL TOO LARGE, DECREASE AS FOLLOWS 002280  

C 002290  

205 DELL=CL(N)/50. 002300  

C 002320  

C PLUME THICKNESS NEGLECTIBLE AT THIS POINT, MOVE ON TO NEXT POINT 002330  

C AFTER SETTING ALL SPECTRAL RADIATION VALUES EQUAL TO ZERO 002340  

C  

IF(DELL.GT..005) GO TO 200  

210 CONTINUE  

C  

DO 211 II=1,NWAV 002390  

RADSTL(II)=0.0  

211 RADSTR(II)=0.0 002400  

GO TO 1990 002450  

C 002460

```

```

*****002470
C      CALCULATE PRESSURE AND TEMPERATURE ALONG LINE-OF-SIGHT AT EACH X,R002480
C                                         002490
C
299 CONTINUE
  WRITE(6,307)DELZ,DELG
307 FORMAT(2(1X,E10.4))
C
DO 300 II=1,NWAV
P(II)=0.0
Q(II)=0.0
TRNS(II)=1.0
300 RADSTR(II)=0.0
C
DO 1000 J=1,N
JCHECK=J
C
DO 301 L=2,NXS
IF(JCHECK.EQ.99) LCHECK=3
NX1=L-1
NX2=L
IF(CX(J).LE.X(L)) GO TO 302
301 CONTINUE
C
302 CONTINUE
DO 303 L=2,25
IF(JCHECK.EQ.99) LCHECK=3
NR2=L
IF(CR(J).LE.R(NX2,L)) GO TO 304
303 CONTINUE
C
304 CONTINUE
DO 305 L=2,25
NR1=L
IF(CR(J).LE.R(NX1,L)) GO TO 306
305 CONTINUE
C
306 CONTINUE
DR1=CR(J)-R(NX1,NR1-1)
DR2=R(NX1,NR1)-R(NX1,NR1-1)
DR3=CR(J)-R(NX2,NR2-1)
DR4=R(NX2,NR2)-R(NX2,NR2-1)
DX1=CX(J)-X(NX1)
DX2=X(NX2)-X(NX1)
C
C     DR1DR2,DR3DR4,DX1DX2 WERE CREATED TO ELIMINATE INDEFINITE FORMS    002760
C
DR1DR2=0.0
DR3DR4=0.0
DX1DX2=0.0
IF(DR2.NE. 0.0) DR1DR2=DR1/DR2
IF(DR4.NE. 0.0) DR3DR4=DR3/DR4
IF(DX2 .NE. 0.0) DX1DX2=DX1/DX2
TX1=T(NX1,NR1-1)+(T(NX1,NR1)-T(NX1,NR1-1))*DR1DR2
TX2=T(NX2,NR2-1)+(T(NX2,NR2)-T(NX2,NR2-1))*DR3DR4
CT=TX1+(TX2-TX1)*DX1DX2
PH20X1=PH20(NX1,NR1-1)+(PH20(NX1,NR1)-PH20(NX1,NR1-1))*DR1DR2
PH20X2=PH20(NX2,NR2-1)+(PH20(NX2,NR2)-PH20(NX2,NR2-1))*DR3DR4
CPH20=PH20X1+(PH20X2-PH20X1)*DX1DX2
C
IF(JCHECK.EQ.99) LCHECK=11
PC02X1=PC02(NX1,NR1-1)+(PC02(NX1,NR1)-PC02(NX1,NR1-1))*DR1DR2
PC02X2=PC02(NX2,NR2-1)+(PC02(NX2,NR2)-PC02(NX2,NR2-1))*DR3DR4
CP02=PC02X1+(PC02X2-PC02X1)*DX1DX2
C
PH2X1=PH2(NX1,NR1-1)+(PH2(NX1,NR1)-PH2(NX1,NR1-1))*DR1DR2
PH2X2=PH2(NX2,NR2-1)+(PH2(NX2,NR2)-PH2(NX2,NR2-1))*DR3DR4

```

$CPN2=PN2X1+(PN2X2-PN2X1)*DX1DX2$
 $P02X1=P02(NX1,NR1-1)+(P02(NX1,NR1)-P02(NX1,NR1-1))*DR1DR2$
 $P02X2=P02(NX2,NR2-1)+(P02(NX2,NR2)-P02(NX2,NR2-1))*DR3DR4$
 $CPO2=P02X1+(P02X2-P02X1)*DX1DX2$

C

$PH2X1=PH2(NX1,NR1)+(PH2(NX1,NR1)-PH2(NX1,NR1-1))*DR1DR2$
 $PH2X2=PH2(NX2,NR2-1)+(PH2(NX2,NR2)-PH2(NX2,NR2-1))*DR3DR4$
 $CPH2=PH2X1+(PH2X2-PH2X1)*DX1DX2$
 $PCO1=PCO(NX1,NR1)+(PCO(NX1,NR1-1)-PCO(NX1,NR1-1))*DR1DR2$
 $PCO2=PCO(NX2,NR2-1)+(PCO(NX2,NR2)-PCO(NX2,NR2-1))*DR3DR4$
 $CPCO=PCO1+(PCO2-PCO1)*DX1DX2$
 $IF(CT.LT.300.) CT=300.$
 $IF(CPH20.LT.0.) CPH20=1.E-20$
 $IF(CPCO2.LT.0.) CPCO2=1.E-20$
 $IF(CPN2.LT.0.) CPN2=1.E-20$
 $IF(CPO2.LT.0.) CPO2=1.E-20$
 $IF(CPH2.LT.0.0)CPH2=1.E-20$
 $IF(CPCO.LT.0.0)CPH2=1.E-20$

002940

AVERAGE TEMPERATURES AND PRESSURES FOR BLOCK

002980

IF(J.GT.1) GO TO 20
 $CPH20L=CPH20$

$CPCO2L=CPCO2$

$CPN2L=CPN2$

$CP02L=CP02$

$CPH2L=CPH2$

$CPCOL=CPCO$

$CTL=CT$

20 CONTINUE

$PCO2AV=(CPCO2+CPCO2L)/2.$
 $PH20AV=(CPH20+CPH20L)/2.$
 $PN2AV=(CPN2+CPN2L)*.5$
 $P02AV=(CP02+CP02L)*.5$
 $PCOAV=(CPCO+CPCOL)/2.$
 $PH2AV=(CPH2+CPH2L)*0.5$

003000

002990

$TAV=(CT+CTL)/2.$

003020

$CPH20L=CPH20$

003030

$CPCO2L=CPCO2$

003040

$CPN2L=CPN2$

$CP02L=CP02$

$CPH2L=CPH2$

$CPCOL=CPCO$

$CTL=CT$

003070

IF(J.EQ.1) GO TO 1000

003080

CALL PARAM(TAV,DEBU,PH20AV,PN2AV,P02AV,PCO2AV,PCOAV,PH2AV)

1002 CONTINUE

DO 2005 II=1,WAV

003100

$TRBLAST(II)=TRNS(II)$

003110

$P(II)=P(II)+IK(II)$

003130

$Q(II)=Q(II)+AKU(II)$

003140

IF (P(II) .EQ. 0.0) P(II)=1.0E-20

003150

IF (Q(II) .EQ. 0.0) Q(II)=1.0E-20

$ARG=-P(II)/ SQRT(1.+.25*P(II)*#2/Q(II))$

003160

$TRNS(II)=EXP(ARG)$

003180

$RAD(II)=(TRBLAST(II)-TRNS(II))*37413./WAV65(II)/(EXP(1.439*WAVNO(II))$

$1.)/TAV)-1.)/3.1415926536$

003200

$RADSTR(II)=RADSTR(II)+RAD(II)$

$SRAD(II)=RAD(II)*10000./WAVNO(II) **2$

$SRADS(II)=SRADS(II)+SRAD(II)$

2005 CONTINUE

1000 CONTINUE

003230

C

```

      WRITE(6,203)
2023 FORMAT(20X,'SPECTRAL RADIATION',//,10X,'WAVNO',10X,'SRADSR')
      DO 2022 II=1,NWAV
2001 STOT = STOT+SRADSR(II)
C   WRITE(6,2024)WAVNO(II),SRADSR(II)
C   ***TAKE OUT WRITE STATEMENT OF SRADSR(II)***_
2024 FORMAT(2(10X,E10.4))
2022 CONTINUE
      CHECK=1,                                     003260
C
C   RAD(II)=RADIATION FOR EACH BLOCK AT EACH WAVELENGTH (WAVLT(II))
C   RAD(II) IN(WATTS/CM2-SR-MICRON)
C   RADX(II) IN (WATTS/CM-SR-MICRON)
C   TRADX IN (WATTS/CM-SR)
C   RADXT(II) IN (WATTS/SR/MICRON)
C   RADSTR(II)=RADIATION AT EACH X AND Z SUMMED OVER EACH PATH LENGTH
C   RADX(II)=RADIATION TOTAL AT EACH X-STATION (SUMMED OVER Z)
C   TRADX=RADX(II) INTEGRATED OVER WAVELENGTH                         003290
C   RADXT(T)=RADX(I) INTEGRATED OVER ALL X-STATIONS                      003300
C   TOTRAD=TTRADX INTEGRATED OVER X                                       003310
C
1980 CONTINUE
1990 IF (M.EQ.1) GO TO 1998                                     003320
      CHECK=2,
      IF(ALPHA.NE.0..AND.ALPHA.NE.180.) GO TO 1996
1981 DO 1982 II=1,NWAV                                         003360
      IZ=II-1
1982 RADXT(II)=RADXT(II)+(RADSTR(II)+RADSTL(II))/2.*PI*DELZ**2*((ZM-1.)003370
      1**2-(ZM-2.)**2)                                              003380
      CHECK=3.
C   THIS CARD HAS BEEN CHANGED AS
C   IF (M.EQ.NZPEX) GO TO 3001
C   IF (M.EQ.NZU) GO TO 3001
      GO TO 1998                                         003410
1996 DO 1997 II=1,NWAV                                         003420
C   INTEGRATED RADIATION IS DOUBLED TO INCLUDE BOTH HALVES OF PLUME 003430
      RADX(II)=RADX(II)+(RADSTR(II)+RADSTL(II))
C   THIS CARD HAS BEEN CHANGED AS
C   IF (M.EQ.NZPEX) RAD(II)=RADX(II)+RADSTR(II)
C   IF (M.EQ.NZU) RAD(II)=RADX(II)+RADSTR(II)
1997 CONTINUE
1998 DO 1999 II=1,NWAV                                         003450
      IZ=II-1
1999 RADSTL(II)=RADSTR(II)                                      003460
2000 CONTINUE
      WRITE(6,2002)STOT
2002 FORMAT(1X,'SUM UP SRADSR AT X STATION',5X,E10.4)
      WRITE(6,1402)
1402 FORMAT(8X,'WAVE NUMBER',5X,'WAVE LENGTH',5X,'RADIATION',
      $ 23X,'WAVE NUMBER',5X,'WAVE LENGTH',6X,'RADIATION')
      WRITE(6,14021)
14021 FORMAT(42X,'/SR/CM/MICRON')
C
      LINCT=0
      DO 12 II=1,NWAV,2
      RADX(II)=RADX(II)*DELZ
      IF(II.GT.NWAV) RADX(II+1)=RADX(II+1)*DELZ
      IF(RADX(II).LT.0.) WRITE(6,18) RADX(II)
18 FORMAT(' NEGATIVE RADIATION IS ILLOGICAL ',E14.4)
      IF((II+1).GT.NWAV) GO TO 15
      WRITE(6,16) WAVNO(II),WAVLT(II),RADX(II),
      $  WAVNO(II+1),WAVLT(II+1),RADX(II+1)
16 FORMAT(5X,F10.0,F16.4,7X,F12.6,20X,F10.0,F16.4,7X,F12.5)
      LINCT=LINCT+1
      IF(LINCT.GT.54) GO TO 12
      LINCT=0

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```

12 CONTINUE
GO TO 17
15 CONTINUE
WRITE(6,16) WAVNO(II),WAVLT(II),RADX(II)
17 CONTINUE
C
1403 FORMAT (1X,F10.0,F16.4,7X,E12.5) 003530
DO 2997 II=2,NWAV 003560
TRADX=TRADX+(WAVLT(II-1)-WAVLT(II))*.
S (RADX(II-1)+RADX(II))*.5
2997 CONTINUE
IF (I.EQ.1) GO TO 2999 003710
CENTER = CENTER+(X(I)-X(I-1))*0.5*(TRADX+TRADXL)*X(I)
TOTRAD=TOTRAD+(TRADX+TRADXL)*0.5*(X(I)-X(I-1)) 003730
STRAD(I)=(TRADX+TRADXL)*0.5*(X(I)-X(I-1))
DO 2998 II=1,NWAV 003740
RADXT(II)=RADXT(II)+((RADX(II)+RADXL(II))/2.
S)*(X(I)-X(I-1))*SC
2998 CONTINUE
WRITE(6,5000) TOTRAD
5000 FORMAT(1H0,24HTOTAL RADIATION EMITTED=E10.4,11H WATTS/STER)
WRITE(6,1402)
DO 5001 II=1,NWAV,2
IF((II+1) .GT. NWAV) GO TO 5002
WRITE(6,16) WAVNO(II),WAVLT(II),RADXT(II),WAVNO(II+1),WAVLT(II+1)
+ ,RADXT(II+1)
5001 CONTINUE
5002 CONTINUE
WRITE(6,16) WAVNO(II),WAVLT(II),RADXT(II)
2999 TRADXL=TRADX 003760
DO 2995 II=1,NWAV 003770
2995 RADXL(II)=RADX(II) 003780
RRN(I)=TRADX 003790
WRITE(6,1492) I,TRADX 003800
1492 FORMAT(/5X,'TOTAL RADIATION FOR STATION',I3,1X,1H=,1X,E12.5,
1      !WATTS/SR/CM!)
C
C CHECK FOR COMBINE RUN. IF SO OUTPUT DATA TO DISC
C
IF(LSWIT .EQ.2) GO TO 22
WRITE(1IDISC) NWAV,I,TRADX
WRITE(1IDISC) (WAVNO(JJ),WAVLT(JJ),RADX(JJ),JJ=1,NWAV)
TDSCT=IDSCT+1
22 CONTINUE
3000 CONTINUE
WRITE(6,1401) 003820
CENTR=0.0 003830
003880
3001 CONTINUE
3003 CENTR=CENTER *SC/TOTRAD 003920
TOTRAD=TOTRAD*SC
IF (TOTRAD .EQ.0.) TOTRAD=1.*10.**(-30)
IF(LSWIT .EQ.1) WRITE(TDSCT) TOTRAD,(RADXT(J),J=1,NWAV)
3004 WRITE(6,1404) (TITLE(J),J=1,8) 003910
WRITE(6,1405)
1405 FORMAT(70X,'STATION RADIATION (W/SR)')
WRITE(6,1407)(X(II),PRN(II),STRAD(II),RCALC(II),II=1,NXS)
1407 FORMAT(215.2,28X,E13.7,20X,S13.7,1X,A1)
C
WRITE(6,8421) 004030
C CALL FINAL PLUT ROUTINE
C
C
8421 FORMAT(1/0X,83H* RADIATION VALUE WAS CALCULATED BASED ON THE SLOPE004040
* BETWEEN THE PREVIOUS TWO POINTS) 004050
SRITE(6,1211) TOTRAD
1211 FORMAT(1H0,24HTOTAL RADIATION EMITTED=E10.4,11H WATTS/STER) 004080

```

```

      IF(LSWIT .EQ.1) WRITE(IDISC) TOTRAD          004090
      WRITE(6,1254) CENTR
      1254 FORMAT(1H0,9HCENTR!D,E11.4,'CM')
C
C   CHECK FOR COMBINED RUN. SENSE SWITCH 4 WILL OVER RIDE THE ISTOP OPTION
C
      IF(LSWIT .EQ.1) GO TO 24
      GO TO (107,103,172),ISTOP                  004280
C
C   DETERMINE IF ANOTHER PASS IS TO BE PERFORMED
C
      24 CONTINUE
      JPASS=JPASS+1
      IDISC=IDISC+1
      KSWIT=1
      IF(JPASS.LT.2) GO TO 47
      KSWIT=2
      JPASS=0
      IDISC=8
C
C   READ H2O AND CO2 RADIATION BACK FROM DISC.
C   SUM AND OUTPUT
C
      END FILE 9
      REWIND 9
      END FILE 8
      REWIND 8
      WRITE(6,27)
      27 FORMAT(1H1,30X,' C O M B I N E D H 2 O C O 2 R A D I A T I O N
      $ S U M S ')
      WRITE(6,28) NWAV1,NWAV2
      28 FORMAT(1H ,/43X,'SUMMED OVER WAVE NUMBERS ',2I6)
C
      DO 25 JK=1,NXS
      READ(8) NWAV,I,CTRADX
      READ(9) NWAV,T,HTRADX
      CALL ERRCHK(2)
      READ(8) (WAVNO(JJ),WAVLT(JJ),CRADX(JJ),JJ=1,NWAV)
      READ(9) (WAVNO(JJ),WAVLT(JJ),HRADX(JJ),JJ=1,NWAV)
      CALL ERRCHK(3)
      DO 26 JJ=1,NWAV
      SRADX(JJ)=HRADX(JJ)+CRADX(JJ)
      26 CONTINUE
      HCOTDX(JK)=HTRADX+CTRADX
C
      WRITE(6,9911)
      WRITE(6,1402)
      DO 29 JJ=1,NWAV,2
      IF((JJ+1).GT.NWAV) GO TO 30
      WRITE(6,16) WAVNO(JJ),WAVLT(JJ),SRADX(JJ),
      S   WAVNO(JJ+1),WAVLT(JJ+1),SRADX(JJ+1)
      29 CONTINUE
      GO TO 31
      30 CONTINUE
      WRITE(6,16) WAVNO(JJ),WAVLT(JJ),SRADX(JJ)
      31 CONTINUE
C
      WRITE(6,32) I,HCTDX(JK)
      32 FORMAT(1H0//10X,' TOTAL COMBINED RADIATION FOR STATION ',I3,
      S   ' IS ',E12.4)
      IF(HCTDX(I)<.00,.0.) GO TO 25
      25 CONTINUE
C
      READ(8) CTRADX,(CRADX(I),J=1,NWAV)
      READ(9) HTRADX,(HRADX(I),J=1,NWAV)
      CALL ERRCHK(4)

```

```

        WRITE(6,9911)
        WRITE(6,34)
34 FORMAT(20X,'WAVNO',10X,'WAVELENGTH',10X,'HCRADXT//')
C PUNCH CARDS FOR ATMOSPHERIC TRANSMISSION PROGRAM
NN=NWAV
DO 33 JJ=1,NWAV
HCRADXT(JJ)=HRADXT(JJ)+CRADXT(JJ)
WRITE(6,35) WAVNO(JJ),WAVLT(JJ),HCRADXT(JJ)
35 FORMAT(17X,E12.4,6X,E12.4,6X,E12.4)
33 CONTINUE
STORAD=HTORAD+CTORAD
WRITE(6,43) (TITLE(J),J=1,8)
43 FORMAT(1H //,' SUMMARIES FOR ',7A10,A2/)
WRITE(6,36) HTORAD
WRITE(6,37) CTORAD
36 FORMAT(10X,'TOTAL H2O RADIATION',E14.4)
37 FORMAT(10X,'TOTAL CO2 RADIATION',E14.4)
WRITE(6,41)
41 FORMAT(29X,-----)
WRITE(6,42) STORAD
42 FORMAT(' TOTAL COMBINED RADIATION',4X,E14.4)

```

C GO TO (107,107,172),ISTOP

107 CONTINUE

STOP

END

SUBROUTINE READER

004420

C
C THIS ROUTINE WILL READ ONE COMPLETE BLOCK OF
C MOTOR DATA PER CALL
C

```

COMMON /INDATA/ X(40),R(40,30),T(40,30),PCO2(40,30),
S PH2O(40,30),PN2(40,30),PO2(40,30),PCO(40,30),PH2(40,30)
COMMON /CTROL/ TITLE(12),KOUNT,NXS,NORUN
COMMON /OTH/ PP,RJ,NRP(40)
COMMON /RASP/RASP
DATA JJ/1/

```

C
READ(5,100) (TITLE(J),J=1,12)
WRITE(6,100) (TITLE(J),J=1,12)
100 FORMAT(12A4)
~~READ(5,101)RASP~~ *This is the radius of the detector. This value does not effect*
~~READ(5,300) PP,RJ,NRP(1)~~
~~RJ=RJ*100.~~
~~WRITE(6,300) PP,RJ,NRP(1)~~
300 FORMAT(2E10.4,14)
KOUNT=0

C
C NOW READ IN THE ENTIRE BLOCK OF MOTOR DATA
C

```

30 CONTINUE
READ(5,400) X(JJ),LAST
X(JJ)=X(JJ)*RJ
WRITE(6,700) X(JJ)
400 FORMAT( E10.3,6X,T1)
700 FORMAT(1X,E10.3)
IF(LAST.EQ.1) GO TO 14
103 FORMAT(15)
READ(5,103)NRP(JJ)
IX=NRP(JJ)
01 12 KK=1,IX
C PUNCH DATA INPUT FOR RADIATION PROGRAM IN
C CHANGE READ (5,500) R(JJ,KK), T(JJ,KK), CARD 10
C THE FOLLOWING ORDER PPCO2, PH2O, PN2, PO2
C +PO2(JJ,KK), PN2(JJ,KK)
C READ(5,500) R(JJ,KK), T(JJ,KK), PCO2(JJ,KK), PH2O(JJ,KK).

```

```

      READ(5,500)R(JJ,KK),T(JJ,KK),PC02(JJ,KK),PH20(JJ,KK)
500 FORMAT(E10.3,20X,3E10.3)
      R(JJ,KK)=R(JJ,KK)*RJ
      PC0(JJ,KK)=1.0E-20
      PH2(JJ,KK)=1.0E-20
      PD2(JJ,KK)=1.0E-20
      PN2(JJ,KK)=1.0E-20
      WRITE(6,600) R(JJ,KK),T(JJ,KK),PN2(JJ,KK),PC0(JJ,KK),PC02(JJ,KK),
+PH20(JJ,KK),PD2(JJ,KK)
600 FORMAT (1X,8E10.4)
12 CONTINUE
      JJ=JJ+1
101 FORMAT(E10.4)
      KOUNT=KOUNT+1
      GO TO 30
14 CONTINUE
      CALL LIMIT(KOUNT,40,1)
      CALL LIMIT(NRP(1),30,2)

C
      RETURN
END
SUBROUTINE LIMIT(IARG,LIM,NPS)
C
IF(IARG.GE.0.AND.IARG.LE.LIM) RETURN
C
      WRITE(6,10) IARG,LIM,NPS
10 FORMAT(' ILLEGAL INPUT PARAMETER ',
$      ' 2110/! CALLED FROM LOCATION ',J8)
C
      STOP
END
SUBROUTINE SWAP
C
COMMON /BANDP/ XKK(254,7),DDINV(254,7),UK(254),AKU(254)
COMMON /RPN1/ NWAV,WAVND(254),WAVLT(254)
C
      DIMENSION XSAVE(254,7),DSAVE(254,7)
C
      DO 10 JJ=1,NWAV
      DO 10 KK=1,7
      XHOLD=XKK(JJ,KK)
      DHOLD=DDINV(JJ,KK)
      XKK(JJ,KK)=XSAVE(JJ,KK)
      DDINV(JJ,KK)=DSAVE(JJ,KK)
      XSAVE(JJ,KK)=XHOLD
      DSAVE(JJ,KK)=DHOLD
10 CONTINUE
C
      NWVOLD=NWAV
      NWAV=NLIST
      NWLST=NVOLD
C
      RETURN
C
ENTRY SAVE1
C
      DO 11 JJ=1,NWAV
      DO 11 KK=1,7
      XSAVE(JJ,KK)=XKK(JJ,KK)
      DSAVE(JJ,KK)=DDINV(JJ,KK)
11 CONTINUE
C
      NVLST=NWAV
C
      RETURN
END

```

SUBROUTINE PRMD

C
COMMON /INDATA/ X(40),R(40,30),T(40,30),PC02(40,30),
\$ PH20(40,30),PH2(40,30),PO2(40,30),PO0(40,30),PH2(40,30)
COMMON /CTRL/ TITLE(12),KOUNT,NXS,NORUN
COMMON /UTH/ PP,RJ,NRP(40)
C
WRITE(6,10)
10 FORMAT(1H1)
WRITE(6,11) NORUN
11 FORMAT(1H ,25X,' M O T C R ',I3//)
WRITE(6,12) KOUNT
12 FORMAT(1H ,9X,' NUMBER OF X STATIONS IS ',I4)
C
DO 14 JJ=1,KOUNT
NRP1=NRP(JJ)
WRITE(6,15) JJ,X(JJ),NRP1
15 FORMAT(1H ,I2,I1)=!,E12.4,
\$ 8X,' NUMBER OF POINTS AT THIS X STATION IS ',I4)
WRITE(6,16)
16 FORMAT(1H ,4X,!R!,13X,!R!,11X,!C02!,10X,
\$ '!H20!,11X,!N2!,11X,!C2!/)
DO 14 KK=1,NRP1
WRITE(6,18) R(JJ,KK),T(JJ,KK),PC02(JJ,KK),
\$ PH20(JJ,KK),PN2(JJ,KK),PO2(JJ,KK)
18 FORMAT(1H ,6(E12.4,1X))
14 CONTINUE

C
RETURN
END

SUBROUTINE BPPRT

C
C THIS ROUTINE PRINTS THE BAND PARAMETERS TOGETHER
C WITH THE ASSOCIATED WAVENUMBERS AND WAVELENGTHS
C

COMMON /BANDP/ XKK(254,7),DDINV(254,7),UK(254),AKU(254)
COMMON /RPN1/ NWAV,WAVNO(254),WAVLT(254)
C
WRITE(6,11)
11 FORMAT(1H1,51X,'B A N D P A R M E T E R S ! //)
LINCT=0
WRITE(6,12)
12 FORMAT(10X,'WAVENUMBER',1X,'WAVELENGTH',20X,
\$ '!XKK AND DDINV PARAMETERS ! /)
C
DO 10 JJ=1,NWAV
WRITE(6,14) WAVNO(JJ),WAVLT(JJ),(XKK(JJ,K),K=1,7)
WRITE(6,14) WAVNO(JJ),WAVLT(JJ),(DDINV(JJ,K),K=1,7)
14 FORMAT(14X,F5.0,4X,F8.1, 7X,7E12.4)
WRITE(6,15)
15 FORMAT(1H)
LINCT=LINCT+3
IF(LINCT.LT.56) GO TO 10
LINCT=0
WRITE(6,16)
16 FORMAT(1H1)
WRITE(6,12)
10 CONTINUE

C
RETURN
END
SUBROUTINE LERITS(CARRAY,HD,RLIM,RUPPER)
C
READ ARRAY(254)
C
RLIM=1,E=10

```

DO 10 JJ=1,NO
IF(ARRAY(JJ).GT.RHIGH) RHIGH=ARRAY(JJ)
10 CONTINUE
C
IUP=ANLOG10(RHIGH)+1
IUP=IUP-3
RUPPER=10.*IUP
RLOW=10.*IUP
C
RETURN
END
SUBROUTINE ERRCHK(ILOC)
RETURN
END
SUBROUTINE PARAM(TAV,XL,PH20,PN2,P02,PCO2,PCO,PH2)
COMMON/CHECK/IREAD
COMMON /BANDP/ XKK(254,7),DDINV(254,7),UK(254),AKU(254)
COMMON /RPN1/ NWA,V,NAVNO(254),WAVLT(254)
COMMON /SWITCH/ KSWIT ,IXXJG
REAL TT(7)
REAL LOG1,LOG2
C
DATA TT(1),TT(2),TT(3),TT(4),TT(5),TT(6),TT(7)/300.,600.,1000.,
$1500.,2000.,2500.,3000./
DATA IDBG/-1/ 004470
T=TAV
IXXJG=IXXJG+ 1
IF (KSWIT.NE.2) GO TO 23
TT(3)=1200.
TT(4)=1500.
TT(5)=1800.
TT(6)=2400.
23 CONTINUE 004480
IDBG=IDBG+1 004490
C CALCULATE OPTICAL DEPTH, U, AND WIDTH, GAMMA
C
T273=273./T
SQT273=SQRT(T273)
IF(KSWIT .EQ.2) GO TO 18
U=T273*PH20*XL
GAMMA=120.+PH20/T+SQT273*(.09*(PH20+PN2)+.04*P02+.12*PCO2
$ +.05*PH2+.1*PCO)
GO TO 19
18 CONTINUE
U=T273*PCO2*XL
GAMMA=.01*PCO2*T273+(.07*PN2+.055*P02+.09*PCO2+PH2*.08
$ +PCO*.06)*SQT273
19 CONTINUE
C
C SELECTION OF TEMPERATURES TO BE USED IN INTERPOLATIONS 004530
C
I=1
DO 30 KK=1,7
KK=KK
TT1=ABS(T-TT(K))
IF (TT1.LE.5-20) GO TO 31
30 CONTINUE
GO TO 32
31 XK=XKK(I,K)
DUV=DDINV(I,K)
GO TO 16
32 CONTINUE
DO 33 J=3,7
K=1
IF(TT(J)-T)< 8,8,9
33 CONTINUE
C
CONTINUE

```

```

IF(T.GE.TT(K-1)) K=K+1
TTK=TT(K)-TT(K-2)
TTK1=TT(K)-TT(K-1)
TTK2=TT(K-1)-TT(K-2)
DIV=TTK*TTK1*TTK2
C
1000 CONTINUE
LOG1=ALOG(XKK(I,K-1)/XKK(I,K-2))
LOG2=ALOG(XKK(I,K)/XKK(I,K-2))
B=((TTK**2)*LOG1-(TTK2**2)*LOG2)/DIV
C=(TTK2*LOG2-TTK*LOG1)/DIV
XK=EXP(ALOG(XKK(I,K-2))+B*(T-TT(K-2))+C*(T-TT(K-2))**2)
C
IF(T.LT.TT(2)) GO TO 14
DLOG=ALOG(DDINV(J,K-2))
DINV=EXP((ALOG(DDINV(I,K-1))-DLOG)*(T-TT(K-2))/(TT(K-1)-TT(K-2))
C +DLOG)
GO TO 16
C
14 CONTINUE
DLOG=ALOG(DDENV(I,1))
DINV=EXP((ALOG(DDINV(I,2))-DLOG)*(T-TT(1))/(TT(2)-TT(1))+DLOG)
C
16 CONTINUE
C CALCULATE PARAMETERS
C
UK(I)=XK*U
AKU(I)=GAMMA*DINV*UK(I)
IF (TT1.LE.E-20) GO TO 17
C
KDEX=K-2
IDEX=K-1
FLOW=.9*AMIN1(XKK(I,KDEX),XKK(I,IDEK))
FHIGH=1.1*AMAX1(XKK(I,KDEX),XKK(I,IDEK))
IF(XK.GE.FLOW.AND.XK.LE.FHIGH) GO TO 17
B=ALOG(XKK(I,K-1)/XKK(I,K-2))/TTK2
XK=EXP(ALOG(XKK(I,K-2))+B*(T-TT(K-2)))
IF(XK.LT.FLOW.OR.XK.GT.FHIGH) WRITE(6,20)
$ WAVNO(I),I,XK,T,U,GAMMA,TTK,R,C,DINV,UK(I),AKU(I)
20 FORMAT(1X,/,! XK STILL OUT OF RANGE AT WAVNO = ',F10.0,' AND I = '
$,1S,/,1X,5E14.4/,1X,5E14.4)
17 CONTINUE
I=I+1
IF (I-NNAV)29,29,15
29 IF (TT1.LE.E-20) GO TO 31
IF(I-NNAV) 1000,1000,15
C
15 CONTINUE
IREAD=IREAD+1
RETURN
END
//GD.FT08F001 DD DSN=HARVE,PLDT8,UNIT=SYSDA,DISP=(NEW,DELETE),
// DCB=(RECFM=VSB,LRECL=100,BLKSIZE=1000,BUFNO=1),
// SPACE=(CYL,(5,5),RLSE)
//GD.FT09F001 DD DSN=HARVF,PLDT9,UNIT=SYSDA,DISP=(NEW,DELETE),
// DCB=(RECFM=VSB,LRECL=100,BLKSIZE=1000,BUFNO=1),
// SPACE=(CYL,(5,5),RLSE)
//GD.SYSTJ 00 *

```

APPENDIX II

INPUT INSTRUCTIONS AND SAMPLE INPUT DATA

APPENDIX II

The computer program input is composed of two parts: the flow field section and the band parameter section. The flow field of a plume can be generated by several computer codes like the NASA MOC code or the AeroChem code [3 and 5]. The flow field program provides the temperature and partial pressure at different locations along the plume. This information is necessary to calculate the optical depth and the radiance of the plume. The band parameter input is composed of the coefficient of absorption part and the line density part. It is tabulated according to the wave number and temperature. Subroutine READER reads in the flow field data and subroutine PARAM calculates the transmissivity τ .

Table A.1 lists all input parameters, format of the read statements, and their meaning, in the same order as encountered in the program for the flow field. Table A.2 lists the input parameters required for the radiation calculations. Tables A.1 and A.2 indicate all input data needed for operation of the code. A listing of a sample input data case is included as Table A.3.

TABLE A.1

<u>Card #</u>	<u>Parameter</u>	<u>Format</u>	<u>Meaning</u>
1	Title	12A4	Title of the Flowfield
2	PP	E10.4	Total Pressure
	RJ	E10.4	Radius of the Nozzle
	NRP(1)	I4	# of Radial Points at the Beginning of the Plume
3	X(JJ)	E10.3	Axial Location
	ILAST	69X,I1	Control value. ILAST = 1 for the end of the flow field
4	NRP(JJ)	I5	# of Radial Points at Axial Location X(JJ)
5	R(JJ,KK)	E10.3	Radial Location at Point (JJ,KK)
	T(JJ,KK)	E10.3	Temperature at Point (JJ,KK)
	PCO ₂ (JJ,KK)	E10.3	Partial Pressure of CO ₂ at Point (JJ,KK)
	PH ₂ O(JJ,KK)	E10.3	Partial Pressure of H ₂ O at Point (JJ,KK)

Repeat card 5 for NRP(JJ) times.

Go back to card 3 for another axial point.

At the end of the flow field, set ILAST = 1. Such value of ILAST will terminate the reader subroutine and return to the MAIN ROUTINE.

TABLE A.2

<u>Card #</u>	<u>Parameter</u>	<u>Format</u>	<u>Meaning</u>
1	NWAV 1	I10	Wave number, lower limit
	NWAV 2	I10	Wave number, upper limit
2	WAV NO(I)	F9.1	Wave number
	XKK(I,J)	7E9.2	Coefficient of absorption K @ 300,600,1200,1500, 1800,2400,3000 °K of wave number WAVNO(I).
3	WAVNO(I)	F9.1	Wave number (must be the same as WAVNO(I) in card 2).
	DDINV(I,J)	7E9.2	Line density at temperature of 300,600,1200,1500,1800, 2400,3000 °K of wave number WAVNO(I).

Repeat card 2 and 3 for N = (NWAV 2 - NWAV 1)/25 +1

4	ALPHA ISTOP	F10.5 I5	Aspect angle of the line of sight ISTOP = 1, will terminate the program after single alpha calculation. ISTOP = 2 will return to this point for another ALPHA calculation. ISTOP = 3, will read in a new set of flow field.
---	----------------	-------------	--

30 June 1978

TABLE A.3

PR1.0X

.1000E 010.3930E 00 12	
0.0	
0.0	0.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.3930E-010.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.7860E-010.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.1179E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.1572E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.1965E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.2358E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.2751E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.3144E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.3537E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.3930E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19	
0.4606E 000.2880E 030.1000E-190.1000E-190.2885E-060.2885E-060.1000E-19	
0.410E 00	
0.0	0.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19
0.3993E-010.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19	
0.7982E-010.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19	
0.1197E 000.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19	
0.1596E 000.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19	
0.1996E 000.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19	
0.2395E 000.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19	
0.2794E 000.2477E 040.1000E-190.1000E-190.7953E-010.2625E 000.1000E-19	
0.3193E 000.2475E 040.1000E-190.1000E-190.7947E-010.2623E 000.1000E-19	
0.3591E 000.2448E 040.1000E-190.1000E-190.7834E-010.2592E 000.1000E-19	
0.3997E 000.2228E 040.1000E-190.1000E-190.6844E-010.2279E 000.1000E-19	
0.4531E 000.1212E 040.1000E-190.1000E-190.2709E-010.8465E-010.1000E-19	
0.5439E 000.5520E 030.1000E-190.1000E-190.6866E-020.2061E-010.1000E-19	
0.6449E 000.3539E 030.1000E-190.1000E-190.1660E-020.4986E-020.1000E-19	
0.7459E 000.3020E 030.1000E-190.1000E-190.3467E-030.1067E-020.1000E-19	
0.8363E 000.2880E 030.1000E-190.1000E-190.2885E-060.2885E-060.1000E-19	
0.814E 00	
0.0	0.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19
0.3997E-010.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19	
0.7994E-010.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19	
0.1199E 000.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19	
0.1599E 000.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19	
0.1999E 000.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19	
0.2398E 000.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19	
0.2798E 000.2486E 040.1000E-190.1000E-190.8207E-010.2607E 000.1000E-19	
0.3197E 000.2473E 040.1000E-190.1000E-190.8185E-010.2594E 000.1000E-19	
0.3596E 000.2412E 040.1000E-190.1000E-190.8035E-010.2512E 000.1000E-19	
0.4016E 000.2173E 040.1000E-190.1000E-190.7212E-010.2137E 000.1000E-19	
0.4559E 000.1479E 040.1000E-190.1000E-190.4261E-010.1220E 000.1000E-19	
0.5357E 000.8166E 030.1000E-190.1000E-190.1581E-010.4558E-010.1000E-19	
0.6308E 000.5012E 030.1000E-190.1000E-190.5747E-020.1687E-010.1000E-19	
0.7318E 000.3739E 030.1000E-190.1000E-190.2240E-020.6671E-020.1000E-19	
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 0.3161E 010.3305E 030.1000E-190.1000E-190.1687E-020.2422E-020.1000E-19
 0.3453E 010.3147E 030.1000E-190.1000E-190.1056E-020.1529E-020.1000E-19
 0.3741E 010.3034E 030.1000E-190.1000E-190.6082E-030.8861E-030.1000E-19
 0.4024E 010.2949E 030.1000E-190.1000E-190.2730E-030.3991E-030.1000E-19
 0.4303E 010.2901E 030.1000E-190.1000E-190.8437E-040.1236E-030.1000E-19
 0.101E 02
 0.0 0.2521E 040.1000E-190.1000E-190.1131E 000.1942E 000.1000E-19
 0.9511E-010.2490E 040.1000E-190.1000E-190.1127E 000.1890E 000.1000E-19
 0.1891E 000.2460E 040.1000E-190.1000E-190.1122E 000.1828E 000.1000E-19
 0.2861E 000.2350E 040.1000E-190.1000E-190.1090E 000.1660E 000.1000E-19
 0.3901E 000.2239E 040.1000E-190.1000E-190.1057E 000.1502E 000.1000E-19
 0.6300E 000.1816E 040.1000E-190.1000E-190.8450E-010.1061E 000.1000E-19
 0.9129E 000.1323E 040.1000E-190.1000E-190.5351E-010.6543E-010.1000E-19
 0.1233E 010.9345E 030.1000E-190.1000E-190.3035E-010.3857E-010.1000E-19
 0.1569E 010.6676E 030.1000E-190.1000E-190.1750E-010.2278E-010.1000E-19
 0.1906E 010.5409E 030.1000E-190.1000E-190.1062E-010.1414E-010.1000E-19
 0.2235E 010.4523E 030.1000E-190.1000E-190.6730E-020.9144E-020.1000E-19
 0.2554E 010.3954E 030.1000E-190.1000E-190.4387E-020.6063E-020.1000E-19
 0.2863E 010.3602E 030.1000E-190.1000E-190.2897E-020.4060E-020.1000E-19
 0.3164E 010.3357E 030.1000E-190.1000E-190.1904E-020.2699E-020.1000E-19
 0.3458E 010.3186E 030.1000E-190.1000E-190.1223E-020.1749E-020.1000E-19
 0.3747E 010.3069E 030.1000E-190.1000E-190.7486E-030.1078E-020.1000E-19
 0.4032E 010.2986E 030.1000E-190.1000E-190.4176E-030.6042E-030.1000E-19
 0.4311E 010.2926E 030.1000E-190.1000E-190.1823E-030.2645E-030.1000E-19
 0.105E 02

0.0	0.2504E	040.1000E-190.1000E-190.1139E	000.1880E	000.1000E-19
0.9432E-010.	2470E	040.1000E-190.1000E-190.1132E	000.1825E	000.1000E-19
0.1914E	000.2441E	040.1000E-190.1000E-190.1126E	000.1770E	000.1000E-19
0.2928E	000.2330E	040.1000E-190.1000E-190.1090E	000.1615E	000.1000E-19
0.3946E	000.2220E	040.1000E-190.1000E-190.1053E	000.1460E	000.1000E-19
0.6398E	000.1801E	040.1000E-190.1000E-190.8391E-010.	1025E	000.1000E-19
0.9208E	000.1341E	040.1000E-190.1000E-190.5492E-010.	6639E-010.	1000E-19
0.1237E	010.9680E	030.1000E-190.1000E-190.3242E-010.	4012E-010.	1000E-19
0.1570E	010.7205E	030.1000E-190.1000E-190.1925E-010.	2440E-010.	1000E-19
0.1905E	010.5682E	030.1000E-190.1000E-190.1195E-010.	1550E-010.	1000E-19
0.2235E	010.4741E	030.1000E-190.1000E-190.7724E-020.	1023E-010.	1000E-19
0.2556E	010.4139E	030.1000E-190.1000E-190.5140E-020.	6924E-020.	1000E-19
0.2867E	010.3737E	030.1000E-190.1000E-190.3473E-020.	4749E-020.	1000E-19
0.3170E	010.3464E	030.1000E-190.1000E-190.2354E-020.	3257E-020.	1000E-19
0.3467E	010.3272E	030.1000E-190.1000E-190.1574E-020.	2200E-020.	1000E-19
0.3757E	010.3134E	030.1000E-190.1000E-190.1015E-020.	1428E-020.	1000E-19
0.4044E	010.3030E	030.1000E-190.1000E-190.5980E-030.	8466E-030.	1000E-19
0.4327E	010.2948E	030.1000E-190.1000E-190.2727E-030.	3872E-030.	1000E-19
0.4606E	010.2901E	030.1000E-190.1000E-190.8526E-040.	1212E-030.	1000E-19
0.110E 02				
0.0	0.2485E	040.1000E-190.1000E-190.1145E	000.1820E	000.1000E-19
0.9528E-010.	2453E	040.1000E-190.1000E-190.1136E	000.1767E	000.1000E-19
0.1938E	000.2421E	040.1000E-190.1000E-190.1128E	000.1714E	000.1000E-19
0.2959E	000.2310E	040.1000E-190.1000E-190.1090E	000.1570E	000.1000E-19
0.3989E	000.2200E	040.1000E-190.1000E-190.1049E	000.1419E	000.1000E-19
0.6351E	000.1809E	040.1000E-190.1000E-190.8426E-010.	1043E	000.1000E-19
0.9169E	000.1333E	040.1000E-190.1000E-190.5430E-010.	6648E-010.	1000E-19
0.1235E	010.9522E	030.1000E-190.1000E-190.3143E-010.	3941E-010.	1000E-19
0.1569E	010.7046E	030.1000E-190.1000E-190.1840E-010.	2363E-010.	1000E-19
0.1906E	010.5549E	030.1000E-190.1000E-190.1129E-010.	1484E-010.	1000E-19
0.2235E	010.4634E	030.1000E-190.1000E-190.7232E-020.	9697E-020.	1000E-19
0.2555E	010.4053E	030.1000E-190.1000E-190.4766E-020.	6502E-020.	1000E-19
0.2865E	010.3669E	030.1000E-190.1000E-190.3186E-020.	4409E-020.	1000E-19
0.3168E	010.3411E	030.1000E-190.1000E-190.2130E-020.	2983E-020.	1000E-19
0.3462E	010.3231E	030.1000E-190.1000E-190.1402E-020.	1982E-020.	1000E-19
0.3752E	010.3103E	030.1000E-190.1000E-190.8881E-030.	1264E-020.	1000E-19
0.4036E	010.3010E	030.1000E-190.1000E-190.5152E-030.	7373E-030.	1000E-19
0.4319E	010.2939E	030.1000E-190.1000E-190.2323E-030.	3335E-030.	1000E-19
0.0				

APPENDIX III

SAMPLE CASE OUTPUT

WAVE NUMBER	WAVE LENGTH	RADIATION W/SR/CM/MICRON	WAVE NUMBER	WAVE LENGTH	RADIATION
2000.	5.0000	0.98137E-03	2025.	4.9383	0.20115E-02
2050.	4.8780	0.39446E-02	2075.	4.8193	0.76348E-02
2100.	4.7619	0.17107E-01	2125.	4.7059	0.29450E-01
2150.	4.6512	0.49690E-01	2175.	4.5977	0.77929E-01
2200.	4.5455	0.11446E 00	2225.	4.4944	0.15660E 00
2250.	4.4444	0.19604E 00	2275.	4.3956	0.22414E 00
2300.	4.3478	0.25091E 00	2325.	4.3011	0.22054E 00
2350.	4.2553	0.17916E 00	2375.	4.2105	0.78166E-01
2400.	4.1667	0.0	2425.	4.1237	0.0
2450.	4.0816	0.0	2475.	4.0404	0.0
2500.	4.0000	0.0			

TOTAL RADIATION EMITTED=0.8259E-01 WATTS/STER

WAVE NUMBER	WAVE LENGTH	RADIATION	WAVE NUMBER	WAVE LENGTH	RADIATION
2000.	5.0000	0.11605E-02	2025.	4.9383	0.23865E-02
2050.	4.8780	0.45994E-02	2075.	4.8193	0.91333E-02
2100.	4.7619	0.19295E-01	2125.	4.7059	0.33014E-01
2150.	4.6512	0.55326E-01	2175.	4.5977	0.85780E-01
2200.	4.5455	0.12457E 00	2225.	4.4944	0.16815E 00
2250.	4.4444	0.20694E 00	2275.	4.3956	0.23378E 00
2300.	4.3478	0.26030E 00	2325.	4.3011	0.22397E 00
2350.	4.2553	0.18512E 00	2375.	4.2105	0.82563E-01
2400.	4.1667	0.0	2425.	4.1237	0.0
2450.	4.0816	0.0	2475.	4.0404	0.0
2500.	4.0000	0.0			

TOTAL RADIATION FOR STATION 4 = 0.78207E-01WATTS/SR/CM

PRI-0X
 STATION NO.= 28 AXIAL DIST. (CM)= 11.035 ASPECT ANGLE (DEG)= 90.0

SUM UP SHADSR AT X STATION 0.0

WAVE NUMBER	WAVE LENGTH	RADIATION	WAVE NUMBER	WAVE LENGTH	RADIATION
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W/SR/CM/MICRON

2000.	5.0000	0.0	2025.	4.9383	0.0
2050.	4.8780	0.0	2075.	4.8193	0.0
2100.	4.7619	0.0	2125.	4.7059	0.0
2150.	4.6512	0.0	2175.	4.5977	0.0
2200.	4.5455	0.0	2225.	4.4944	0.0
2250.	4.4444	0.0	2275.	4.3956	0.0
2300.	4.3473	0.0	2325.	4.3011	0.0
2350.	4.2553	0.0	2375.	4.2105	0.0
2400.	4.1667	0.0	2425.	4.1237	0.0
2450.	4.0816	0.0	2475.	4.0404	0.0
2500.	4.0000	0.0			

TOTAL RADIATION EMITTED=0.1663E 01 WATTS/SIER

WAVE NUMBER	WAVE LENGTH	RADIATION	WAVE NUMBER	WAVE LENGTH	RADIATION
2000.	5.0000	0.10455E-01	2025.	4.9383	0.21407E-01
2050.	4.8780	0.41942E-01	2075.	4.8193	0.51130E-01
2100.	4.7619	0.26070E 00	2125.	4.7059	0.47098E 00
2150.	4.6512	0.82891E 00	2175.	4.5977	0.13972E 01
2200.	4.5455	0.21785E 01	2225.	4.4944	0.31900E 01
2250.	4.4444	0.42430E 01	2275.	4.3956	0.50712E 01
2300.	4.3473	0.56955E 01	2325.	4.3011	0.51826E 01
2350.	4.2553	0.39990E 01	2375.	4.2105	0.16854E 01
2400.	4.1667	0.0	2425.	4.1237	0.0
2450.	4.0815	0.0	2475.	4.0404	0.0
2500.	4.0000	0.0			

TOTAL RADIATION FOR STATION 28 = 0.0 WATTS/SR/CM

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LAPP GAS DYNAMICS - ARMY RADIATION CODE

PRI.DX	W/SR/cm	STATION RADIATION (W/SR)
0.0	0.5837254E-01	0.4032077E-00
0.41	0.6468284E-01	0.2522010E-01
0.81	0.7061887E-01	0.2730462E-01
1.22	0.7820714E-01	0.3006312E-01
1.61	0.8649528E-01	0.3262293E-01
1.97	0.9388489E-01	0.3232552E-01
2.37	0.1016850E 00	0.3915963E-01
2.78	0.1099497E 00	0.4262582E-01
3.18	0.1185645E 00	0.4656448E-01
3.58	0.1266638E 00	0.4842820E-01
3.98	0.1350878E 00	0.5225714E-01
4.39	0.1435786E 00	0.5804345E-01
4.82	0.1524286E 00	0.6281853E-01
5.22	0.1602380E 00	0.6266773E-01
5.63	0.1665467E 00	0.6678170E-01
6.03	0.1720112E 00	0.6852239E-01
6.44	0.1782773E 00	0.7039573E-01
6.85	0.1849118E 00	0.7493514E-01
6.94	0.1851284E 00	0.1672365E-01
7.41	0.1922510E 00	0.8824474E-01
7.85	0.1983629E 00	0.8673418E-01
8.30	0.2038116E 00	0.9088081E-01
8.76	0.2088537E 00	0.9325266E-01
9.21	0.2134077E 00	0.9542012E-01
9.66	0.2167625E 00	0.9805322E-01
10.12	0.2199754E 00	0.9955025E-01
10.58	0.2249740E 00	0.1014207E 00
11.04	0.0	0.5172247E-01

* RADIATION VALUE WAS CALCULATED BASED ON THE SLOPE BETWEEN THE PREVIOUS TWO POINTS

TOTAL RADIATION EMITTED=0.1663E 01 WATTS/SIER

CENTROID= 0.6655E 01CM

NWAV1	2000	2500
WAV	2000.0	0.78E-02
WAV	2000.0	0.95E-01
WAV	2025.0	0.54E-02
WAV	2025.0	0.12E-00
WAV	2050.0	0.38E-02
WAV	2050.0	0.11E-00
WAV	2075.0	0.26E-02
WAV	2075.0	0.12E-00
WAV	2100.0	0.18E-02
WAV	2100.0	0.14E-00
WAV	2125.0	0.13E-02
WAV	2125.0	0.20E-00
WAV	2150.0	0.88E-03
WAV	2150.0	0.22E-00
WAV	2175.0	0.62E-02
WAV	2175.0	0.23E-00
WAV	2200.0	0.48E-03
WAV	2200.0	0.23E-00
WAV	2225.0	0.41E-03
WAV	2225.0	0.24E-00
WAV	2250.0	0.32E-03
WAV	2250.0	0.25E-00
WAV	2275.0	0.23E-03
WAV	2275.0	0.25E-00
WAV	2300.0	0.20E-03
WAV	2300.0	0.27E-00
WAV	2325.0	0.15E-03
WAV	2325.0	0.27E-00
WAV	2350.0	0.10E-03
WAV	2350.0	0.28E-00
WAV	2375.0	0.85E-04
WAV	2375.0	0.24E-00
WAV	2400.0	0.76E-04
WAV	2400.0	0.30E-00
WAV	2425.0	0.62E-04
WAV	2425.0	0.31E-00
WAV	2450.0	0.48E-04
WAV	2450.0	0.32E-00
WAV	2475.0	0.37E-04
WAV	2475.0	0.33E-00
WAV	2500.0	0.36E-04
WAV	2500.0	0.34E-00